

## NONLOCAL ELLIPTIC EQUATIONS IN BOUNDED DOMAINS: A SURVEY

XAVIER ROS-OTON

**Abstract:** In this paper we survey some results on the Dirichlet problem

$$\begin{cases} Lu = f & \text{in } \Omega \\ u = g & \text{in } \mathbb{R}^n \setminus \Omega \end{cases}$$

for nonlocal operators of the form

$$Lu(x) = \text{PV} \int_{\mathbb{R}^n} \{u(x) - u(x+y)\} K(y) dy.$$

We start from the very basics, proving existence of solutions, maximum principles, and constructing some useful barriers. Then, we focus on the regularity properties of solutions, both in the interior and on the boundary of the domain.

In order to include some natural operators  $L$  in the regularity theory, we do not assume any regularity on the kernels. This leads to some interesting features that are purely nonlocal, in the sense that they have no analogue for local equations.

We hope that this survey will be useful for both novel and more experienced researchers in the field.

**2010 Mathematics Subject Classification:** 47G20, 60G52, 35B65.

**Key words:** Integro-differential equations, bounded domains, regularity.

### CONTENTS

1. Introduction	4
2. Motivation and some preliminaries	6
3. Existence of solutions	9
4. Comparison principle	11
5. Barriers and $L^\infty$ bounds	13
6. Interior regularity	16
7. Boundary regularity	18
8. Further interior regularity	20
References	21

## 1. Introduction

The aim of this paper is to survey some results on Dirichlet problems of the form

$$(1.1) \quad \begin{cases} Lu = f & \text{in } \Omega \\ u = g & \text{in } \mathbb{R}^n \setminus \Omega. \end{cases}$$

Here,  $\Omega$  is any bounded domain in  $\mathbb{R}^n$ , and  $L$  is an elliptic integro-differential operator of the form

$$(1.2) \quad Lu(x) = \text{PV} \int_{\mathbb{R}^n} \{u(x) - u(x+y)\} K(y) dy.$$

The function  $K(y) \geq 0$  is the kernel of the operator<sup>1</sup>, and we assume

$$K(y) = K(-y) \quad \text{and} \quad \int_{\mathbb{R}^n} \min\{|y|^2, 1\} K(y) dy < \infty.$$

Integro-differential problems of the form (1.1) arise naturally in the study of stochastic processes with jumps, and have been widely studied both in Probability and in Analysis and PDEs. We refer to the reader to the following works on:

- Existence of solutions: Felsinger–Kassmann–Voigt [25], Leonori–Peral–Primo–Soria [39], Hoh–Jacob [31], and Barles–Imbert [5].
- Interior regularity of solutions: Bass–Levin [6], Kassmann [33], Caffarelli–Silvestre [13, 14, 15], Barles–Chasseigne–Imbert [4], Kassmann–Mimica [34, 35], Schwab–Silvestre [53], and Serra [54].
- Boundary regularity of solutions: Bogdan [8], the author and Serra [46, 47, 48], Grubb [29, 30], Chen–Song [16], Barles–Chasseigne–Imbert [3], Bogdan–Grzywny–Ryznar [9], and Bogdan–Kumagai–Kwaśnicki [10].
- Other qualitative properties of solutions: Birkner–López–Mimbela–Wakolbinger [7], Dipierro–Savin–Valdinoci [20], Abatangelo [1], and Kulczycki [36].

In applications, this type of nonlocal problems appear in models involving “anomalous diffusions” (in which the underlying stochastic process is not given by Brownian motion), or in the presence of long-range interactions or forces. In particular, they appear in Physics [21, 59,

---

<sup>1</sup>The typical example is  $K(y) = c_{n,s}|y|^{-n-2s}$ , which corresponds to  $L = (-\Delta)^s$ , the fractional Laplacian. As  $s \uparrow 1$ , it converges to the Laplacian  $-\Delta$ .

For equations with  $x$ -dependence, the kernel is a function of two variables,  $K(x, y)$ . Here, we assume that the operator  $L$  is translation invariant, and thus that the kernel  $K$  does not depend on  $x$ .

**43, 27, 23, 38**], Finance [**52, 40, 2**], Fluid dynamics [**18, 17**], Ecology [**42, 32, 45**], or Image processing [**28**].

The aim of this paper is to survey some results on existence, regularity, and qualitative properties of solutions to (1.1). More precisely, we will start from the very basics (showing existence and boundedness of solutions), to then focus on the regularity of solutions, both in the interior and on the boundary of the domain  $\Omega$ .

In order to include some natural operators  $L$  in the regularity theory, we do *not* assume any regularity on the kernel  $K(y)$ . As we will see, there is an interesting relation between the regularity properties of solutions and the regularity of the kernels  $K(y)$ . This is a purely nonlocal issue, in the sense that has no analogue in second order equations, as explained next.

For linear elliptic second order equations of the form

$$(1.3) \quad \begin{cases} -\sum_{ij} a_{ij} \partial_{ij} u = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega, \end{cases}$$

the interior regularity properties of  $u$  depend only on the regularity of  $f$ . This is because, after an affine change of variables, this equation is just  $-\Delta u = f$ .

The nonlocal analogue of (1.3) is (1.1). For this problem, the interior regularity of solutions depends on the regularity of  $f$  – as in (1.3) –, but it also depends on the regularity of  $K(y)$  in the  $y$ -variable. Furthermore, if the kernel  $K$  is not regular, then the interior regularity of  $u$  will in addition depend on the regularity of  $g$ , on the boundary regularity of  $u$ , and even on the shape of  $\Omega$ , as we will see later on.

Hence, for nonlocal equations, the class of linear and translation invariant operators is much richer, and already presents several interesting features. The same type of issues appear when considering nonlinear equations, or operators with  $x$ -dependence. However, for the clarity of presentation, we will consider only linear and translation invariant equations.

In most of the regularity results we will focus on two classes of kernels: we will assume  $s \in (0, 1)$  and either that

$$(1.4) \quad \frac{\lambda}{|y|^{n+2s}} \leq K(y) \leq \frac{\Lambda}{|y|^{n+2s}}, \quad 0 < \lambda \leq \Lambda;$$

or that

$$(1.5) \quad K(y) = \frac{a(y/|y|)}{|y|^{n+2s}}, \quad a \in L^1(S^{n-1}), \quad a \geq 0.$$

In fact, in order to include operators like  $L = (-\partial_{x_1 x_1}^2)^s + \dots + (-\partial_{x_n x_n}^2)^s$ , we will allow  $a$  in (1.5) to be any nonnegative measure on  $S^{n-1}$ . The only necessary condition on the measure  $a$  is that it is not supported in a hyperplane – so that the operator  $L$  is not degenerate. A quantitative way to state this nondegeneracy condition is that, for some positive constants  $\lambda$  and  $\Lambda$ , one has  $\int_{S^{n-1}} da \leq \Lambda$  and  $\int_{S^{n-1}} |\nu \cdot \theta|^{2s} da(\theta) \geq \lambda$  for all  $\nu \in S^{n-1}$ ; see [48].

The paper is organized as follows. In Section 2 we briefly explain the probabilistic interpretation of (1.1). In Section 3 we show existence of weak solutions. In Section 4 we prove the maximum principle for such solutions. In Section 5 we construct some useful barriers and give an  $L^\infty$  estimate. In Section 6 we discuss the interior regularity properties for the classes of kernels (1.4) and (1.5). In Section 7 we see what is the boundary regularity of solutions. Finally, in Section 8 we come back to the interior regularity of solutions to (1.1).

## 2. Motivation and some preliminaries

**2.1. Lévy processes.** Integro-differential equations of the form (1.1) arise naturally in the study of stochastic processes with jumps, and more precisely in Lévy processes. A Lévy process is a stochastic process with independent and stationary increments. Informally speaking, it represents the random motion of a particle whose successive displacements are independent and statistically identical over different time intervals of the same length.

These processes extend the concept of Brownian motion, and were introduced a few years after Wiener gave the precise definition of the Brownian motion [41]. Essentially, Lévy processes are obtained when one relaxes the assumption of continuity of paths (which gives the Brownian motion) by the weaker assumption of stochastic continuity.

By the Lévy–Khintchine Formula, the infinitesimal generator of any Lévy processes is an operator of the form<sup>2</sup>

$$\begin{aligned} -Lu(x) &= \sum_{i,j} a_{ij} \partial_{ij} u + \sum_j b_j \partial_j u \\ &\quad + \int_{\mathbb{R}^n} \{u(x+y) - u(x) - y \cdot \nabla u(x) \chi_{B_1}(y)\} d\nu(y), \end{aligned}$$

where  $\nu$  is the Lévy measure, and satisfies  $\int_{\mathbb{R}^n} \min\{1, |y|^2\} d\nu(y) < \infty$ .

---

<sup>2</sup>We denote the infinitesimal generator  $-L$  in order to be consistent with our notation in (1.2). We always take  $L$  to be positive definite, so it is the analogue of  $-\Delta$ .

When the process has no diffusion or drift part, this operator takes the form

$$-Lu(x) = \int_{\mathbb{R}^n} \{u(x+y) - u(x) - y \cdot \nabla u(x) \chi_{B_1}(y)\} d\nu(y).$$

Furthermore, if one assumes the process to be symmetric, and the Lévy measure to be absolutely continuous, then  $L$  can be written as (1.2) or, equivalently,

$$Lu(x) = \frac{1}{2} \int_{\mathbb{R}^n} \{2u(x) - u(x+y) - u(x-y)\} K(y) dy,$$

with  $K(y) = K(-y)$ . We will use this last expression for  $L$  throughout the paper.

As an example, let  $\Omega \subset \mathbb{R}^n$  be any bounded domain, and let us consider a Lévy process  $X_t$ ,  $t \geq 0$ , starting at  $x \in \Omega$ . Let  $u(x)$  be the expected first exit time, i.e., the expected time  $\mathbb{E}[\tau]$ , where  $\tau = \inf\{t > 0 : X_t \notin \Omega\}$  is the first time at which the particle exits the domain. Then,  $u(x)$  solves

$$\begin{cases} Lu = 1 & \text{in } \Omega \\ u = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \end{cases}$$

where  $-L$  is the infinitesimal generator of  $X_t$ .

Recall that, when  $X_t$  is a Brownian motion, then  $L$  is the Laplace operator  $-\Delta$ . In the context of integro-differential equations, Lévy processes plays the same role that Brownian motion plays in the theory of second order equations.

Another simple example is given by the following. Let us also consider a bounded domain  $\Omega$ , a process  $X_t$  starting at  $x \in \Omega$ , and the first time  $\tau$  at which the particle exits the domain. Assume now that we have a payoff function  $g: \mathbb{R}^n \setminus \Omega \rightarrow \mathbb{R}$ , so that when the process  $X_t$  exits  $\Omega$  we get a payoff  $g(X_\tau)$ . Then, the expected payoff  $u(x) := \mathbb{E}[g(X_\tau)]$  solves the problem

$$\begin{cases} Lu = 0 & \text{in } \Omega \\ u = g & \text{in } \mathbb{R}^n \setminus \Omega. \end{cases}$$

The Dirichlet problem (1.1) arises when considering at the same time a running cost  $f$  and a final payoff  $g$ .

**2.2. Kernels with compact support.** It is important to remark that when the kernel  $K(y)$  has compact support in a ball  $B_\delta$  (for some  $\delta > 0$ ), then the Dirichlet problem is

$$\begin{cases} Lu = f & \text{in } \Omega \\ u = g & \text{in } (\Omega + B_\delta) \setminus \Omega. \end{cases}$$

This means that the exterior condition  $g$  has to be posed only in a neighborhood of  $\partial\Omega$  and not in the whole  $\mathbb{R}^n \setminus \Omega$ . This turns out to be convenient in some applications; see for example [21, 57, 59, 11].

However, from the analytical point of view, it is mainly the singularity of the kernel  $K(y)$  at the origin that plays a role in the regularity properties of solutions; see for example Section 14 in [13].

**2.3. Stable processes.** A special class of Lévy processes are the so-called stable processes. These are the processes which satisfy self-similarity properties, and they are also the ones appearing in the Generalized Central Limit Theorem; see [51].

The infinitesimal generators of these processes are given by (1.2)–(1.5). Note that the structural condition (1.5) on the kernel  $K$  is equivalent to saying that the Lévy measure is *homogeneous*. This is also equivalent to the fact that the operator  $L$  is scale invariant.

A very natural stable process is the radially symmetric one. This means that  $K(y) = c|y|^{-n-2s}$  and hence, up to a multiplicative constant,

$$(2.1) \quad L = (-\Delta)^s.$$

Another very natural stable process is the one obtained by taking independent stable processes in each coordinate. That is, we consider  $X_t = (X_t^{(1)}, \dots, X_t^{(n)})$ , where  $X_t^{(i)}$  are 1-dimensional i.i.d. symmetric stable processes. The generator of this process  $X_t$  will be

$$(2.2) \quad L = (-\partial_{x_1 x_1}^2)^s + \dots + (-\partial_{x_n x_n}^2)^s,$$

and it corresponds to (1.5) with the measure  $a$  being  $2n$  delta functions on  $S^{n-1}$ . For example, when  $n = 2$ , one has  $a = \delta_{(0,1)} + \delta_{(0,-1)} + \delta_{(1,0)} + \delta_{(-1,0)}$ .

While the operators (2.1) and (2.2) may look similar (they are the same when  $s = 1$ ), they have quite different regularity properties, as we will see in the next sections. For example, while solutions to  $(-\Delta)^s u = 0$  in  $\Omega$  are  $C^\infty$  inside the domain, this may not be the case for (2.2).

Note also that the Fourier symbols of (2.1) and (2.2) are  $|\xi|^{2s}$  and  $|\xi_1|^{2s} + \dots + |\xi_n|^{2s}$ , respectively. Again, the first one is  $C^\infty$  outside the origin, while the second one is just  $C^{2s}$ . In general, the Fourier symbol of any stable process of the form (1.5) is

$$(2.3) \quad A(\xi) = c \int_{S^{n-1}} |\xi \cdot \theta|^{2s} a(\theta) d\theta.$$

The symbol  $A(\xi)$  is in general only  $C^{2s}$  outside the origin, but it would be  $C^\infty$  outside the origin whenever the function  $a$  is  $C^\infty$  on  $S^{n-1}$ .

### 3. Existence of solutions

In this section we explain briefly how to prove existence of weak solutions to (1.1) in a simple case. We hope that this will be useful for students and/or for readers which are not familiar with nonlocal operators. The more experienced reader should go to Felsinger–Kassmann–Voigt [25], where this is done in a much more general setting.

As we will see, the existence (and uniqueness) of weak solutions to (1.1)–(1.2) follows from the Riesz representation theorem once one has the appropriate ingredients.

The energy functional associated to the problem (1.1) is

$$(3.1) \quad \mathcal{E}(u) = \frac{1}{4} \iint_{\mathbb{R}^{2n} \setminus (\Omega^c \times \Omega^c)} (u(x) - u(z))^2 K(z - x) dx dz - \int_{\Omega} fu.$$

The minimizer of  $\mathcal{E}$  among all functions with  $u = g$  in  $\mathbb{R}^n \setminus \Omega$  will be the unique weak solution of (1.1).

Notice that  $\mathcal{E}(u)$  is defined for all regular enough functions  $u$  which are bounded at infinity. When  $g \equiv 0$ , then making the change of variables  $y = x + z$  the functional can be written as

$$(3.2) \quad \mathcal{E}(u) = \frac{1}{4} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} (u(x) - u(x + y))^2 K(y) dx dy - \int_{\Omega} fu.$$

When  $g$  is not zero, the term  $\iint_{\Omega^c \times \Omega^c} |g(x) - g(z)|^2 K(z - x) dx dz$  could be infinite, and this is why in general one has to take (3.1).

For simplicity, we will show existence of solutions for the case  $g \equiv 0$ , and hence we can think on the energy functional (3.2).

Let  $H_K(\mathbb{R}^n)$  be the space of functions  $u \in L^2(\mathbb{R}^n)$  satisfying

$$(3.3) \quad [u]_{H_K}^2 := \frac{1}{2} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} (u(x) - u(x + y))^2 K(y) dx dy < \infty,$$

and let

$$X = \{u \in H_K(\mathbb{R}^n) : u \equiv 0 \text{ in } \mathbb{R}^n \setminus \Omega\}.$$

Essentially, the only assumption which is needed in order to prove existence of solutions is the Poincaré inequality

$$(3.4) \quad \int_{\Omega} u^2 \leq C \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |u(x) - u(x+y)|^2 K(y) dx dy$$

for functions  $u \equiv 0$  in  $\mathbb{R}^n \setminus \Omega$ .

When the kernel  $K$  satisfies (1.4), then the seminorm in (3.3) is equivalent to

$$(3.5) \quad [u]_{H^s}^2 := \frac{1}{2} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \frac{|u(x) - u(x+y)|^2}{|y|^{n+2s}} dx dy.$$

In this case, the Poincaré inequality (3.4) follows easily from the fractional Sobolev inequality<sup>3</sup> in  $\mathbb{R}^n$  and Hölder's inequality in  $\Omega$ .

Once one has the Poincaré inequality (3.4), then it follows that the space  $X$  is a Hilbert space with the scalar product

$$(v, w)_K := \frac{1}{2} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} (v(x) - v(x+y))(w(x) - w(x+y))K(y) dx dy.$$

Then, the weak formulation of (1.1) is just

$$(3.6) \quad (u, \varphi)_K = \int_{\Omega} f \varphi \quad \text{for all } \varphi \in X,$$

and the existence and uniqueness of weak solution follows immediately from the Riesz representation theorem.

For more general classes of kernels  $K(y)$ , one has to show the Poincaré inequality (3.4) (see Lemma 2.7 in [25]), and then the existence of solutions follows by the same argument above.

---

<sup>3</sup>Different proofs of the fractional Sobolev inequality can be found in [58], [44], and [19]. The proof given in [44], which is due to Brezis, is really nice and simple.

Finally, notice that when  $u$  is regular enough then for all  $\varphi \in X$  we have

$$\begin{aligned}
 (u, \varphi) &= \frac{1}{2} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} (u(x) - u(z))(\varphi(x) - \varphi(z))K(x - z) dx dz \\
 &= \frac{1}{2} \text{PV} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} (u(x) - u(z))\varphi(x)K(x - z) dx dz \\
 &\quad + \frac{1}{2} \text{PV} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} (u(z) - u(x))\varphi(z)K(x - z) dx dz \\
 &= \frac{1}{2} \int_{\mathbb{R}^n} Lu(x) \varphi(x) dx + \frac{1}{2} \int_{\mathbb{R}^n} Lu(z)\varphi(z) dz \\
 &= \int_{\Omega} Lu \varphi,
 \end{aligned}$$

where we used that  $K(y) = K(-y)$  and that  $\varphi \equiv 0$  in  $\mathbb{R}^n \setminus \Omega$ .

This means that, when  $u$  is regular enough, the weak formulation (3.6) reads as

$$\int_{\Omega} Lu \varphi = \int_{\Omega} f \varphi \quad \text{for all } \varphi \in X,$$

and thus it is equivalent to  $Lu = f$  in  $\Omega$ .

*Remark 3.1.* We showed in this section how to prove the existence of *weak* solutions to (1.1) by variational methods. On the other hand, an alternative approach is to use Perron's method to establish the existence of *viscosity* solutions; see [3, 13]. As we will see, when the right hand side  $f$  is Hölder continuous then all solutions are classical solutions (in the sense that the operator  $L$  can be evaluated pointwise), and thus the notions of weak and viscosity solutions coincide; see [56] for more details.

#### 4. Comparison principle

In this section we show the maximum principle and the comparison principle for weak solutions to (1.1).

As in the classical case of the Laplacian  $-\Delta$ , the maximum principle essentially relies on the fact that  $Lu(x_0) \geq 0$  whenever  $u$  has a maximum at  $x_0$ . In case of local equations, this is true for any local maximum, while in nonlocal equations this is only true when  $u$  has a global maximum at  $x_0$ . In this case, the inequality  $Lu(x_0) \geq 0$  follows simply from the expression

$$Lu(x_0) = \frac{1}{2} \int_{\mathbb{R}^n} (2u(x_0) - u(x_0 + y) - u(x_0 - y))K(y) dy$$

and the fact that  $u(x_0) \geq u(x_0 \pm y)$  for any  $y \in \mathbb{R}^n$  (which holds whenever  $u$  has a global maximum at  $x_0$ ).

Notice also that when  $K(y) > 0$  in  $\mathbb{R}^n$ , this argument already yields a strong maximum principle, since one has  $Lu(x_0) > 0$  unless  $u \equiv \text{cst}$ .

The previous considerations work when  $u$  is regular enough, so that  $Lu$  can be evaluated pointwise. In case of weak solutions  $u$  to (1.1), the proof of the maximum principle goes as follows. For simplicity, we do it in the case  $K(y) > 0$ , but the same could be done for more general kernels  $K(y) \geq 0$ .

**Proposition 4.1.** *Let  $L$  be any operator of the form (1.2), with  $K(y) > 0$  in  $\mathbb{R}^n$ . Let  $u$  be any weak solution to (1.1), with  $f \geq 0$  in  $\Omega$  and  $g \geq 0$  in  $\mathbb{R}^n \setminus \Omega$ . Then,  $u \geq 0$  in  $\Omega$ .*

*Proof:* Recall that  $u$  is a weak solution of (1.1) if  $u = g$  in  $\mathbb{R}^n \setminus \Omega$  and

$$(4.1) \quad \iint_{\mathbb{R}^{2n} \setminus (\Omega^c \times \Omega^c)} (u(x) - u(z))(\varphi(x) - \varphi(z))K(z-x) dx dz = \int_{\Omega} f\varphi$$

for all  $\varphi \in H_K(\mathbb{R}^n)$  with  $\varphi \equiv 0$  in  $\mathbb{R}^n \setminus \Omega$ .

Write  $u = u^+ - u^-$  in  $\Omega$ , where  $u^+ = \max\{u, 0\}\chi_{\Omega}$  and  $u^- = \max\{-u, 0\}\chi_{\Omega}$ . We will take  $\varphi = u^-$ , assume that  $u^-$  is not identically zero, and argue by contradiction.

Indeed, since  $f \geq 0$  and  $\varphi \geq 0$ , then we clearly have that

$$(4.2) \quad \int_{\Omega} f\varphi \geq 0.$$

On the other hand, we have that

$$\begin{aligned} & \iint_{\mathbb{R}^{2n} \setminus (\Omega^c \times \Omega^c)} (u(x) - u(z))(\varphi(x) - \varphi(z))K(z-x) dx dz \\ &= \int_{\Omega} \int_{\Omega} (u(x) - u(z))(u^-(x) - u^-(z))K(z-x) dx dz \\ & \quad + 2 \int_{\Omega} dx \int_{\Omega^c} (u(x) - g(z))u^-(x)K(z-x) dz. \end{aligned}$$

Moreover,  $(u^+(x) - u^+(z))(u^-(x) - u^-(z)) \leq 0$ , and thus

$$\begin{aligned} & \int_{\Omega} \int_{\Omega} (u(x) - u(z))(u^-(x) - u^-(z))K(z-x) dx dz \\ & \leq - \int_{\Omega} \int_{\Omega} (u^-(x) - u^-(z))^2 K(z-x) dx dz < 0. \end{aligned}$$

Note that the last strict inequality holds since we are assuming that  $u^-$  is not identically zero, and since  $u^-$  can not be constant (this would mean that  $u$  equals to a negative constant inside  $\Omega$ , which cannot be if  $g \geq 0$  and  $f \geq 0$ ). Also, since  $g \geq 0$  then

$$\int_{\Omega} dx \int_{\Omega^c} (u(x) - g(z)) u^-(x) K(z - x) dz \leq 0.$$

Therefore, we have shown that

$$\iint_{\mathbb{R}^{2n} \setminus (\Omega^c \times \Omega^c)} (u(x) - u(z)) (\varphi(x) - \varphi(z)) K(z - x) dx dz < 0,$$

and this contradicts (4.1)–(4.2).  $\square$

Of course, once we have the maximum principle, the comparison principle follows immediately.

**Corollary 4.2.** *Let  $L$  be any operator of the form (1.2), with  $K(y) > 0$  in  $\mathbb{R}^n$ . Let  $u_1$  and  $u_2$  be weak solutions to*

$$\begin{cases} Lu_1 = f_1 & \text{in } \Omega \\ u_1 = g_1 & \text{in } \mathbb{R}^n \setminus \Omega \end{cases} \quad \text{and} \quad \begin{cases} Lu_2 = f_2 & \text{in } \Omega \\ u_2 = g_2 & \text{in } \mathbb{R}^n \setminus \Omega. \end{cases}$$

*Assume that  $f_1 \geq f_2$  and  $g_1 \geq g_2$ . Then,  $u_1 \geq u_2$ .*

*Proof:* Just apply Proposition 4.1 to  $u = u_1 - u_2$ .  $\square$

This allows us to use barriers, which in turn can be used to show that solutions  $u$  to (1.1) belong to  $L^\infty(\Omega)$  whenever  $f$  and  $g$  are bounded. This is what we do in the next section.

## 5. Barriers and $L^\infty$ bounds

We provide in this section an  $L^\infty$  estimate of the form

$$(5.1) \quad \|u\|_{L^\infty(\Omega)} \leq \|g\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + C \|f\|_{L^\infty(\Omega)}$$

for solutions to (1.1). The reader should go to [39] for the case of right hand sides  $f \in L^p(\Omega)$ .

To do it, we assume first  $K(y) > 0$  in  $\mathbb{R}^n$ , as in the previous section. In this case, the construction of a barrier is quite simple.

**Lemma 5.1.** *Let  $L$  be an operator of the form (1.2), with  $K(y) > 0$  in  $\mathbb{R}^n$ . Then, there exists a function  $w \in C_c^\infty(\mathbb{R}^n)$  such that*

$$\begin{cases} Lw \geq 1 & \text{in } \Omega \\ w \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega \\ w \leq C & \text{in } \Omega. \end{cases}$$

*The constant  $C$  depends only on the kernel  $K$  and  $\text{diam}(\Omega)$ .*

Once we have this barrier, the  $L^\infty$  bound (5.1) follows from the comparison principle, as shown next.

**Corollary 5.2.** *Let  $L$  be any operator of the form (1.2), with  $K(y) > 0$  in  $\mathbb{R}^n$ . Let  $u$  be any weak solution of (1.1). Then,*

$$\|u\|_{L^\infty(\Omega)} \leq \|g\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + C\|f\|_{L^\infty(\Omega)},$$

where  $C$  is the constant in Lemma 5.1.

*Proof:* Let  $v(x) = \|g\|_{L^\infty} + \|f\|_{L^\infty} w(x)$ , where  $w$  is given by Lemma 5.1. Then, we clearly have  $Lv \leq Lv$  in  $\Omega$ , and  $g \leq v$  in  $\mathbb{R}^n \setminus \Omega$ .

Thus, by the comparison principle, we have  $u \leq v$  in  $\Omega$ . In particular,  $u \leq \|g\|_{L^\infty} + C\|f\|_{L^\infty}$  in  $\Omega$ .

Applying the same argument to  $(-u)$ , we find that  $-u \leq \|g\|_{L^\infty} + C\|f\|_{L^\infty}$ , and hence the result follows.  $\square$

We now construct the barrier.

*Proof of Lemma 5.1:* Let  $B_R$  be any large enough ball such that  $\Omega \Subset B_R$ , and let  $\eta \in C_c^\infty(B_R)$  be such that

$$(5.2) \quad 0 \leq \eta \leq 1 \quad \text{in } \mathbb{R}^n, \quad \eta \equiv 1 \quad \text{in } \Omega.$$

Then, for each  $x \in \Omega$  we have  $\eta(x) = \max_{\mathbb{R}^n} \eta$ , and thus

$$2\eta(x) - \eta(x+y) - \eta(x-y) \geq \eta(x) - \eta(x+y) \geq 0.$$

Hence, writing  $z = x + y$ , we have

$$L\eta(x) \geq \int_{\mathbb{R}^n} \{\eta(x) - \eta(z)\} K(x-z) dz \geq \int_{\mathbb{R}^n \setminus B_R} K(x-z) dz,$$

where we have used that  $\eta(x) - \eta(z) = 1$  for  $z$  in  $\mathbb{R}^n \setminus B_R$ . Now, we notice that

$$\int_{\mathbb{R}^n \setminus B_R} K(x-z) dz = \int_{\mathbb{R}^n \setminus (x+B_R)} K(y) dy \geq \int_{\mathbb{R}^n \setminus B_{2R}} K(y) dy = c > 0$$

for some positive constant  $c$ .

Hence, we have  $L\eta \geq c > 0$  in  $\Omega$ . Taking  $w = \frac{1}{c}\eta$ , we will have that

$$\begin{cases} Lw \geq 1 & \text{in } \Omega \\ w \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega \\ w \leq \frac{1}{c} & \text{in } \Omega, \end{cases}$$

as desired.  $\square$

Note that the previous proof works not only for all kernels satisfying  $K(y) > 0$  in  $\mathbb{R}^n$ , but also for all kernels  $K$  satisfying

$$(5.3) \quad \int_{\mathbb{R}^n \setminus B_R} K(y) dy > 0$$

for every ball  $B_R$ . In particular, the above construction works for all stable operators (1.5).

Hence, we see that the desired barrier is quite easy to construct, and does not even depend much on the class of kernels we are considering. Essentially, thanks to the nonlocal character of the operator, any function  $\eta$  as in (5.2) works.

*Remark 5.3.* For operators  $L$  that do not satisfy (5.3), one can still manage to prove an analogous result. Indeed, if (5.3) does not hold, then  $K = 0$  a.e. outside a ball  $B_R$ . This means that  $K$  have compact support. In this case, one can take any function  $\eta \geq 0$  which is strictly concave in a large ball  $B_M$  and zero outside  $B_M$ . Then, if  $M$  is large enough (so that  $\Omega + \text{supp}(K) \subset B_M$  and hence  $2\eta(x) - \eta(x+y) - \eta(x-y) > 0$  therein), one will have  $L\eta \geq c > 0$  in  $\Omega$ .

Finally, to end this section, we give an important explicit solution for the class of stable operators (1.5).

**Lemma 5.4.** *Let  $L$  be any stable operator of the form (1.2)–(1.5). Then, the function<sup>4</sup>*

$$u_0(x) := (1 - |x|^2)_+^s$$

solves

$$\begin{cases} Lu_0 = c & \text{in } B_1 \\ u_0 = 0 & \text{in } \mathbb{R}^n \setminus B_1 \end{cases}$$

for some positive constant  $c > 0$ .

This explicit solution will be used in Section 7 to construct a subsolution and establish a Hopf Lemma for this class of operators.

Lemma 5.4 was first established for  $(-\Delta)^s$  in dimension  $n = 1$  by Gettoor [26], and requires quite fine computations; see also the work of Dyda [22].

Once this is established in dimension  $n = 1$ , the result in dimension  $n$  and for general stable operators  $L$  follows by just writing

$$Lu_0(x) = \frac{1}{4} \int_{S^{n-1}} a(\theta) \left( \int_{\mathbb{R}} \frac{2u_0(x) - u_0(x + \tau\theta) - u_0(x - \tau\theta)}{|\tau|^{1+2s}} d\tau \right) d\theta,$$

---

<sup>4</sup>Here  $z_+$  denotes the positive part of the number  $z$ , i.e.,  $z_+ = \max\{z, 0\}$ .

and using the result in dimension 1 for each direction  $\theta$ . It is important to notice that the 1D functions  $\tau \mapsto u_0(x + \tau\theta)$  are exactly a rescaled version of  $(1 - |\tau|^2)_+^s$ , which solves the equation in dimension  $n = 1$ . Using this, one gets that  $Lu_0 = c$  in  $B_1$ , with  $c = c_s \int_{S^{n-1}} a(\theta) d\theta$ , and  $c_s$  depending only on  $s$ .

## 6. Interior regularity

For second order equations, the classical Schauder estimate for  $-\Delta u = f$  in  $B_1$  establishes that

$$(6.1) \quad \|u\|_{C^{2+\alpha}(B_{1/2})} \leq C (\|f\|_{C^\alpha(B_1)} + \|u\|_{L^\infty(B_1)})$$

whenever  $\alpha > 0$  is not an integer. Thus, it immediately follows from this estimate that solutions to the Dirichlet problem

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases}$$

are  $C^{2+\alpha}$  inside  $\Omega$  whenever  $f \in C^\alpha$  and  $g$  is bounded.

For nonlocal equations of order  $2s$ , one would expect a similar estimate, in which the norm  $\|u\|_{C^{2s+\alpha}(B_{1/2})}$  is controlled by  $\|f\|_{C^\alpha(B_1)}$  for solutions to  $Lu = f$  in  $B_1$ . It turns out that, due to the nonlocality of the equation, the norm  $\|u\|_{L^\infty(B_1)}$  in (6.1) has to be replaced by a global norm of  $u$ , i.e., a norm that controls  $u$  in the whole  $\mathbb{R}^n$ .

**6.1. Regular kernels.** For the fractional Laplacian  $L = (-\Delta)^s$  – which corresponds to  $K(y) = c|y|^{-n-2s}$  in (1.2) –, this estimate reads as

$$(6.2) \quad \|u\|_{C^{2s+\alpha}(B_{1/2})} \leq C (\|f\|_{C^\alpha(B_1)} + \|u\|_{L^\infty(\mathbb{R}^n)}),$$

and holds whenever  $\alpha + 2s$  is not an integer; see for example [37, 46].

The same estimate (6.2) holds for more general operators of the form (1.2)–(1.4) or (1.2)–(1.5) under the extra assumption that the kernels  $K(y)$  are  $C^\alpha$  outside the origin; see [54, Corollary 1.2] and [48, Corollary 3.5] for more details.

This means that, in this case of regular kernels, solutions to (1.1) are  $C^{2s+\alpha}$  inside  $\Omega$  whenever  $f \in C^\alpha$  and  $g$  is bounded.

However, for general operators (1.2)–(1.4) or (1.2)–(1.5) (with no further regularity assumption on the kernel  $K$ ), the estimate (6.2) is not true anymore, and one needs a stronger norm of  $u$  in the right hand side, as explained next.

**6.2. Non-regular kernels.** When the kernels  $K(y)$  are not regular outside the origin, one has the following estimate.

**Theorem 6.1** ([54, 48]). *Let  $L$  be any operator of the form (1.2) with kernel  $K$  of the form either (1.4) or (1.5). Let  $\alpha > 0$  be such that  $\alpha + 2s$  is not an integer, and let  $u \in L^\infty(\mathbb{R}^n)$  be any weak solution to  $Lu = f$  in  $B_1$ . Then,*

$$(6.3) \quad \|u\|_{C^{2s+\alpha}(B_{1/2})} \leq C (\|f\|_{C^\alpha(B_1)} + \|u\|_{C^\alpha(\mathbb{R}^n)}),$$

for some constant  $C$  that depends only on  $n$ ,  $s$ , and  $\lambda$ , and  $\Lambda$ .

It is important to remark that the previous estimate is valid also in case  $\alpha = 0$  (in which the  $C^\alpha$  norm has to be replaced by the  $L^\infty$  norm); see Theorem 1.1 in [48] for more details.

The proof of this estimate uses a refined version of the blow-up and compactness argument first introduced by Serra in [55]. The estimate (6.3) was established in [54] for the class of kernels (1.4)<sup>5</sup>, and in [48] for the class (1.5).

With no further regularity assumption on the kernels  $K$ , the estimate (6.3) is sharp, in the sense that the norm  $\|u\|_{C^\alpha(\mathbb{R}^n)}$  can not be replaced by a weaker one. More precisely, one has the following.

**Proposition 6.2** ([54, 48]). *Let  $s \in (0, 1)$ ,  $\alpha \in (0, s]$ . Then, for any small  $\epsilon > 0$  there exists an operator of the form (1.2)–(1.5), and a solution  $u$  to  $Lu = 0$  in  $B_1$  such that  $u \in C^{\alpha-\epsilon}(\mathbb{R}^n)$ , but  $u \notin C^{2s+\alpha}(B_{1/2})$ . Moreover, the same happens for the operators (1.2)–(1.4).*

Thus, even when  $f \equiv 0$ , solutions  $u$  to

$$\begin{cases} Lu = 0 & \text{in } B_1 \\ u = g & \text{in } \mathbb{R}^n \setminus B_1 \end{cases}$$

are in general no better than  $C^{2s}(B_{1/2})$  if  $g$  is not better than  $L^\infty(\mathbb{R}^n)$ . This means that the interior regularity of solutions to (1.1) depends on the regularity of  $g$  in  $\mathbb{R}^n \setminus \Omega$  when the kernels are not regular.

Finally notice that, even in case that  $g$  is  $C^\infty$  in all of  $\mathbb{R}^n \setminus \Omega$ , this is not enough to deduce that  $u$  is regular inside  $\Omega$ ! Indeed, because of (6.3), one has to control the term  $\|u\|_{C^\alpha(\mathbb{R}^n)}$  in order to have a  $C^{2s+\alpha}$  estimate inside  $\Omega$ . For this, it is not enough to have  $\|g\|_{C^\alpha(\mathbb{R}^n \setminus \Omega)} \leq C$ , but one also needs to control the regularity of  $u$  across  $\partial\Omega$ . In other words, the boundary regularity of  $u$  is needed. We will come back to this in Section 8.

---

<sup>5</sup>The estimate in [54] is a much stronger result, which holds for fully nonlinear equations and also for  $x$ -dependent equations.

## 7. Boundary regularity

In this section we study the boundary regularity of solutions to

$$(7.1) \quad \begin{cases} Lu = f & \text{in } \Omega \\ u = 0 & \text{in } \mathbb{R}^n \setminus \Omega. \end{cases}$$

We will first look at the optimal Hölder regularity of solutions  $u$  near  $\partial\Omega$ , to then see more fine results for the class of kernels (1.5).

*Remark 7.1.* Once (7.1) is well understood, the boundary regularity of solutions to (1.1) follows from the results for (7.1), at least when the exterior data  $g$  in (1.1) is regular enough.

Indeed, assume that  $u$  is a solution to (1.1) and that  $g \in C^{2s+\gamma}(\mathbb{R}^n \setminus \Omega)$ . We may extend  $g$  to a function in  $\mathbb{R}^n$  satisfying  $g \in C^{2s+\gamma}(\mathbb{R}^n)$ , and consider  $\tilde{u} = u - g$ . Then  $\tilde{u}$  satisfies  $\tilde{u} \equiv 0$  in  $\mathbb{R}^n \setminus \Omega$ , and  $L\tilde{u} = f - Lg =: \tilde{f}$  in  $\Omega$ . Hence,  $\tilde{u}$  solves (7.1) with  $u$  and  $f$  replaced by  $\tilde{u}$  and  $\tilde{f}$ , respectively.

**7.1. Hölder regularity for  $u$ .** As we saw in Section 5, given any stable operator (1.2)–(1.5), the function

$$u_0(x) = (1 - |x|^2)_+^s$$

is an explicit solution to (7.1) in  $\Omega = B_1$ ; see Lemma 5.4.

The function  $u_0$  belongs to  $C^s(\overline{B_1})$ , but

$$u_0 \notin C^{s+\epsilon}(\overline{B_1}) \quad \text{for any } \epsilon > 0.$$

This means that, even in the simplest case  $K(y) = |y|^{-n-2s}$ , one can not expect solutions to be better than  $C^s(\overline{\Omega})$ .

For the class of kernels (1.5), using the explicit solution  $u_0$  and similar barriers, it is possible to show that, when  $\Omega$  is  $C^{1,1}$ , solutions  $u$  satisfy

$$|u| \leq Cd^s \quad \text{in } \Omega, \quad d(x) = \text{dist}(x, \mathbb{R}^n \setminus \Omega).$$

Combining the previous bound with the interior estimate (6.3) (with  $\alpha = 0$ ), one gets the following.

**Proposition 7.2** (Optimal Hölder regularity, [48]). *Let  $L$  be any operator of the form (1.2)–(1.5), and  $\Omega$  any bounded  $C^{1,1}$  domain. Let  $f \in L^\infty(\Omega)$ , and  $u$  be the weak solution of (7.1). Then,*

$$\|u\|_{C^s(\overline{\Omega})} \leq C \|f\|_{L^\infty(\Omega)}$$

for some constant  $C$  that depends only on  $n$ ,  $s$ ,  $\Omega$ ,  $\Lambda$ , and  $\lambda$ .

Furthermore, using similar barriers one can show:

**Lemma 7.3** (Hopf's Lemma). *Let  $L$  be any operator of the form (1.2)–(1.5), and  $\Omega$  any bounded  $C^{1,1}$  domain. Let  $u$  be any weak solution to (7.1), with  $f \geq 0$ . Then, either*

$$u \geq cd^s \quad \text{in } \Omega \quad \text{for some } c > 0$$

or  $u \equiv 0$  in  $\Omega$ .

Thus, the boundary regularity of solutions depends essentially on the construction of suitable barriers, and the behavior of these barriers near  $\partial\Omega$  depends on the class of kernels under consideration.

When the kernels  $K(y)$  are of the form (1.5), we have seen that solutions behave like  $d^s$  near the boundary  $\partial\Omega$ . However, for the class of kernels (1.4), the only barriers that one can construct behave like  $d^{\alpha_0}$  near the boundary, for some  $0 < \alpha_0 < s$ . This means that for this class of operators one can only prove  $u \in C^{\alpha_0}(\bar{\Omega})$ .

**7.2. Regularity of  $u/d^s$ .** We have seen that, when

$$(7.2) \quad K(y) = \frac{a(y/|y|)}{|y|^{n+2s}},$$

then all solutions  $u$  behave like  $d^s$  near the boundary. For this class of kernels, much more can be said about the regularity of solutions near  $\partial\Omega$ .

Indeed, the quotient  $u/d^s$  is not only bounded, but it is also Hölder continuous up to the boundary.

This regularity of  $u/d^s$  yields in particular the existence of the limit

$$\frac{u}{d^s}(z) := \lim_{\Omega \ni x \rightarrow z} \frac{u(x)}{d^s(x)}$$

for all  $z \in \partial\Omega$ . This function  $u/d^s$  on  $\partial\Omega$  plays sometimes the role that the normal derivative  $\partial u/\partial\nu$  plays in second order equations. For example, it appears in overdetermined problems [24], integration by parts formulas [49], and free boundary problems [12].

The first proof of this result was given in [46] for the case  $K(y) = |y|^{-n-2s}$ , and more recently the result has been improved by Grubb in [29, 30] and by the author and Serra in [47, 48]. These results may be summarized as follows.

**Theorem 7.4** ([48, 47, 30, 29]). *Let  $s \in (0, 1)$ ,  $L$  be any operator of the form (1.2)–(1.5),  $\Omega$  be any bounded domain, and  $u$  be any solution to (7.1), with  $f \in L^\infty(\Omega)$ .*

*Then, depending on the regularity of the function  $a$  in (7.2), we have the following.*

(i) If  $a$  is any measure and  $\Omega$  is  $C^{1,1}$ , then

$$\|u/d^s\|_{C^{s-\epsilon}(\overline{\Omega})} \leq C\|f\|_{L^\infty(\overline{\Omega})}.$$

(ii) If  $a \in C^{1,\gamma}(S^{n-1})$  and  $\Omega$  is  $C^{2,\gamma}$  for some small  $\gamma > 0$ , then

$$\|u/d^s\|_{C^{s+\gamma}(\overline{\Omega})} \leq C\|f\|_{C^\gamma(\overline{\Omega})}$$

whenever  $s + \gamma$  is not an integer and  $f \in C^\gamma(\overline{\Omega})$ .

(iii) If  $a \in C^\infty(S^{n-1})$ ,  $\Omega$  is  $C^\infty$ , and  $f \in C^\gamma(\overline{\Omega})$ , then

$$\|u/d^s\|_{C^{s+\gamma}(\overline{\Omega})} \leq C\|f\|_{C^\gamma(\overline{\Omega})}$$

for all  $\gamma \in (0, \infty)$  such that  $s + \gamma$  is not an integer.

Part (i) corresponds to Theorem 1.2 in [48], part (ii) was established in [47] in the more general context of fully nonlinear equations, and part (iii) was established in [29, 30] for all pseudodifferential operators satisfying the  $\mu$ -transmission property. (In the excepted cases of (iii), more information is given in [29] in terms of Hölder–Zygmund spaces  $C_*^k$ , when  $s + \gamma$  is an integer.)

As in the interior regularity estimates, the regularity of the kernel  $K$  affects the regularity of the solution  $u$ . In this direction, even if one assumes that  $\Omega$  is  $C^\infty$  and  $f \in C^\infty(\overline{\Omega})$ , the result in part (i) can not be improved if the kernels are singular; see [48].

On the other hand, notice that these results yield  $(u - g)/d^s \in C^{s+\gamma}(\overline{\Omega})$  for solutions  $u$  to (1.1); see Remark 7.1. This gives a description of  $u$  near  $\partial\Omega$  up to order  $2s + \gamma$ . For example, if  $s + \gamma < 1$ , this means that there exists a function  $b \in C^\gamma(\partial\Omega)$  such that

$$|u(x) - g(x) - b(z)d^s(x)| \leq C|x - z|^{2s+\gamma} \quad \text{for } z \in \partial\Omega, \quad x \in B_1(z).$$

When  $s + \gamma \in (1, 2)$  one has a similar expansion with an additional term of order  $d^{s+1}$ , and more terms appear for higher values of  $\gamma$ .

## 8. Further interior regularity

As we saw in the previous section, solutions  $u$  to (7.1) are  $C^s$  or  $C^{\alpha_0}$  up to the boundary when  $K(y)$  is of the form (1.4) or (1.5). Moreover, the same happens for solutions to (1.1) whenever the exterior data  $g$  is regular enough.

Hence, this means that solutions  $u$  will be globally Hölder continuous, and therefore we may apply the estimate

$$(8.1) \quad \|u\|_{C^{\alpha+2s}(B_{1/2})} \leq C(\|f\|_{C^\alpha(\overline{\Omega})} + \|u\|_{C^\alpha(\mathbb{R}^n)})$$

to get that solutions  $u$  are  $C^{2s+\alpha}$  inside  $\Omega$  for some  $\alpha > 0$  – more precisely,  $\alpha = s$  in case (1.5), and  $\alpha = \alpha_0$  in case (1.4). This is enough to ensure that, when  $f \in C^\alpha$ , any weak solution is a classical solution, in the sense that the operator  $L$  can be evaluated pointwise.

The natural question then is: are solutions more regular than this?

For the class of kernels (1.5) the interior estimate (8.1) does not give more than  $u \in C_{\text{loc}}^{3s}(\Omega)$ , and we know that (8.1) is sharp. Still, if one considers the solution  $u$  to

$$(8.2) \quad \begin{cases} Lu = 1 & \text{in } \Omega \\ u = 0 & \text{in } \mathbb{R}^n \setminus \Omega \end{cases}$$

in a  $C^\infty$  domain  $\Omega$ , one a priori does not know if solutions are more regular than  $C^{3s}$ . This question has been answered by the author and Valdinoci in [50]:

**Theorem 8.1** ([50]). *Let  $L$  be any operator of the form (1.2)–(1.5), and  $\Omega$  be any bounded domain. Then,*

- (i) *If  $a$  is any measure and  $\Omega$  is convex, then  $u \in C_{\text{loc}}^{1+3s-\epsilon}(\Omega)$  for all  $\epsilon > 0$ .*
- (ii) *If  $a \in L^\infty(S^{n-1})$  and  $\Omega$  is  $C^{1,1}$ , then  $u \in C_{\text{loc}}^{1+3s-\epsilon}(\Omega)$  for all  $\epsilon > 0$ .*
- (iii) *There is a (nonconvex)  $C^\infty$  domain  $\Omega$ , and an operator  $L$  of the form (1.2)–(1.5), for which the solution  $u$  to (8.2) does not belong to  $C_{\text{loc}}^{3s+\epsilon}(\Omega)$  for any  $\epsilon > 0$ .*

Thus, solutions have different regularity properties depending on the shape of  $\Omega$ . Indeed, for operators  $L$  with singular kernels, the convexity of the domain gives one more order of differentiability on the solution. Also, notice that when  $a \in L^\infty$ , then we gain interior regularity on the solution by assuming that the domain  $\Omega$  is  $C^{1,1}$ .

For kernels  $K$  in the class (1.4), one would expect solutions  $u$  to be  $C^{1+2s+\alpha_0-\epsilon}$  in the interior of  $C^{1,1}$  domains.

## References

- [1] N. ABATANGELO, Large  $s$ -harmonic functions and boundary blow-up solutions for the fractional Laplacian, *Discrete Contin. Dyn. Syst.* **35**(12) (2015), 5555–5607. DOI: 10.3934/dcds.2015.35.5555.
- [2] V. AKGIRAY AND G. G. BOOTH, The stable-law model of stock returns, *J. Bus. Econ. Stat.* **6**(1) (1988), 55–57. DOI: 10.1080/07350015.1988.10509636.

- [3] G. BARLES, E. CHASSEIGNE, AND C. IMBERT, On the Dirichlet problem for second-order elliptic integro-differential equations, *Indiana Univ. Math. J.* **57(1)** (2008), 213–246. DOI: 10.1512/iumj.2008.57.3315.
- [4] G. BARLES, E. CHASSEIGNE, AND C. IMBERT, Hölder continuity of solutions of second-order non-linear elliptic integro-differential equations, *J. Eur. Math. Soc. (JEMS)* **13(1)** (2011), 1–26. DOI: 10.4171/JEMS/242.
- [5] G. BARLES AND C. IMBERT, Second-order elliptic integro-differential equations: viscosity solutions’ theory revisited, *Ann. Inst. H. Poincaré Anal. Non Linéaire* **25(3)** (2008), 567–585. DOI: 10.1016/j.anihpc.2007.02.007.
- [6] R. F. BASS AND D. A. LEVIN, Harnack inequalities for jump processes, *Potential Anal.* **17(4)** (2002), 375–388. DOI: 10.1023/A:1016378210944.
- [7] M. BIRKNER, J. A. LÓPEZ-MIMBELA, AND A. WAKOLBINGER, Comparison results and steady states for the Fujita equation with fractional Laplacian, *Ann. Inst. H. Poincaré Anal. Non Linéaire* **22(1)** (2005), 83–97. DOI: 10.1016/j.anihpc.2004.05.002.
- [8] K. BOGDAN, The boundary Harnack principle for the fractional Laplacian, *Studia Math.* **123(1)** (1997), 43–80.
- [9] K. BOGDAN, T. GRZYWNY, AND M. RYZNAR, Barriers, exit time and survival probability for unimodal Lévy processes, *Probab. Theory Related Fields* **162(1)** (2015), 155–198. DOI: 10.1007/s00440-014-0568-6.
- [10] K. BOGDAN, T. KUMAGAI, AND M. KWAŚNICKI, Boundary Harnack inequality for Markov processes with jumps, *Trans. Amer. Math. Soc.* **367(1)** (2015), 477–517. DOI: 10.1090/S0002-9947-2014-06127-8.
- [11] L. CAFFARELLI, S. PATRIZI, AND V. QUITALO, On a long range segregation model, Submitted (2014).
- [12] L. A. CAFFARELLI, J.-M. ROQUEJOFFRE, AND Y. SIRE, Variational problems with free boundaries for the fractional Laplacian, *J. Eur. Math. Soc. (JEMS)* **12(5)** (2010), 1151–1179. DOI: 10.4171/JEMS/226.
- [13] L. CAFFARELLI AND L. SILVESTRE, Regularity theory for fully nonlinear integro-differential equations, *Comm. Pure Appl. Math.* **62(5)** (2009), 597–638. DOI: 10.1002/cpa.20274.
- [14] L. CAFFARELLI AND L. SILVESTRE, Regularity results for nonlocal equations by approximation, *Arch. Ration. Mech. Anal.* **200(1)** (2011), 59–88. DOI: 10.1007/s00205-010-0336-4.

- [15] L. CAFFARELLI AND L. SILVESTRE, The Evans–Krylov theorem for nonlocal fully nonlinear equations, *Ann. of Math. (2)* **174**(2) (2011), 1163–1187. DOI: 10.4007/annals.2011.174.2.9.
- [16] Z.-Q. CHEN AND R. SONG, Estimates on Green functions and Poisson kernels for symmetric stable processes, *Math. Ann.* **312**(3) (1998), 465–501. DOI: 10.1007/s002080050232.
- [17] P. CONSTANTIN, Euler equations, Navier–Stokes equations and turbulence, in: “*Mathematical foundation of turbulent viscous flows*”, Lecture Notes in Math. **1871**, Springer, Berlin, 2006, pp. 1–43. DOI: 10.1007/11545989\_1.
- [18] A.-L. DALIBARD AND D. GÉRARD-VARET, On shape optimization problems involving the fractional Laplacian, *ESAIM Control Optim. Calc. Var.* **19**(4) (2013), 976–1013. DOI: 10.1051/cocv/2012041.
- [19] E. DI NEZZA, G. PALATUCCI, AND E. VALDINOCI, Hitchhiker’s guide to the fractional Sobolev spaces, *Bull. Sci. Math.* **136**(5) (2012), 521–573. DOI: 10.1016/j.bulsci.2011.12.004.
- [20] S. DIPIERRO, O. SAVIN, AND E. VALDINOCI, All functions are locally  $s$ -harmonic up to a small error, *J. Eur. Math. Soc. (JEMMS)* (to appear).
- [21] Q. DU, M. GUNZBURGER, R. B. LEHOUCQ, AND K. ZHOU, Analysis and approximation of nonlocal diffusion problems with volume constraints, *SIAM Rev.* **54**(4) (2012), 667–696. DOI: 10.1137/110833294.
- [22] B. DYDA, Fractional calculus for power functions and eigenvalues of the fractional Laplacian, *Fract. Calc. Appl. Anal.* **15**(4) (2012), 536–555. DOI: 10.2478/s13540-012-0038-8.
- [23] A. C. ERINGEN, “*Nonlocal continuum field theories*”, Springer-Verlag, New York, 2002.
- [24] M. M. FALL AND S. JAROHS, Overdetermined problems with fractional Laplacian, *ESAIM Control Optim. Calc. Var.* (to appear).
- [25] M. FELSINGER, M. KASSMANN, AND P. VOIGT, The Dirichlet problem for nonlocal operators, *Math. Z.* **279**(3) (2015), 779–809. DOI: 10.1007/s00209-014-1394-3.
- [26] R. K. GETOOR, First passage times for symmetric stable processes in space, *Trans. Amer. Math. Soc.* **101** (1961), 75–90. DOI: 10.1090/S0002-9947-1961-0137148-5.
- [27] G. GIACOMIN AND J. L. LEBOWITZ, Phase segregation dynamics in particle systems with long range interactions. I. Macroscopic limits, *J. Statist. Phys.* **87**(1) (1997), 37–61. DOI: 10.1007/BF02181479.

- [28] G. GILBOA AND S. OSHER, Nonlocal operators with applications to image processing, *Multiscale Model. Simul.* **7(3)** (2008), 1005–1028. DOI: 10.1137/070698592.
- [29] G. GRUBB, Local and nonlocal boundary conditions for  $\mu$ -transmission and fractional elliptic pseudodifferential operators, *Anal. PDE* **7(7)** (2014), 1649–1682. DOI: 10.2140/apde.2014.7.1649.
- [30] G. GRUBB, Fractional Laplacians on domains, a development of Hörmander’s theory of  $\mu$ -transmission pseudodifferential operators, *Adv. Math.* **268** (2015), 478–528. DOI: 10.1016/j.aim.2014.09.018.
- [31] W. HOH AND N. JACOB, On the Dirichlet problem for pseudodifferential operators generating Feller semigroups, *J. Funct. Anal.* **137(1)** (1996), 19–48. DOI: 10.1006/jfan.1996.0039.
- [32] N. E. HUMPHRIES ET AL. Environmental context explains Lévy and Brownian movement patterns of marine predators, *Nature* **465** (2010), 1066–1069. DOI: 10.1038/nature09116.
- [33] M. KASSMANN, A priori estimates for integro-differential operators with measurable kernels, *Calc. Var. Partial Differential Equations* **34(1)** (2009), 1–21. DOI: 10.1007/s00526-008-0173-6.
- [34] M. KASSMANN AND A. MIMICA, Intrinsic scaling properties for nonlocal operators, Preprint (2013). arXiv:1310.5371.
- [35] M. KASSMANN AND A. MIMICA, Intrinsic scaling properties for nonlocal operators II, Preprint (2015). arXiv:1412.7566.
- [36] T. KULCZYCKI, On concavity of solution of Dirichlet problem for the equation  $(-\Delta)^{1/2}\varphi = 1$  in a convex planar region, Preprint (2014). arXiv:1405.3846.
- [37] N. S. LANDKOF, “*Foundations of modern potential theory*”, Translated from the Russian by A. P. Doohovskoy, Die Grundlehren der mathematischen Wissenschaften **180**, Springer-Verlag, New York-Heidelberg, 1972.
- [38] N. LASKIN, Fractional quantum mechanics and Lévy path integrals, *Phys. Lett. A* **268(4–6)** (2000), 298–305. DOI: 10.1016/S0375-9601(00)00201-2.
- [39] T. LEONORI, I. PERAL, A. PRIMO, AND F. SORIA, Basic estimates for solutions of a class of nonlocal elliptic and parabolic equations, *Discrete Contin. Dyn. Syst.* **35(12)** (2015), 6031–6068. DOI: 10.3934/dcds.2015.35.6031.
- [40] S. Z. LEVENDORSKIĬ, Pricing of the American put under Lévy processes, *Int. J. Theor. Appl. Finance* **7(3)** (2004), 303–335. DOI: 10.1142/S0219024904002463.
- [41] P. LÉVY, “*Théorie de l’addition des variables aléatoires*”, Monographies des probabilités **1**, Gauthier-Villars, Paris, 1937.

- [42] A. MASSACCESI AND E. VALDINOCI, Is a nonlocal diffusion strategy convenient for biological populations in competition?, Preprint (2015). [arXiv:1503.01629](#).
- [43] R. METZLER AND J. KLAFTER, The random walk's guide to anomalous diffusion: a fractional dynamics approach, *Phys. Rep.* **339(1)** (2000), 1–77. DOI: 10.1016/S0370-1573(00)00070-3.
- [44] A. C. PONCE, “*Elliptic PDEs, Measures and Capacities*”, Tracts in Mathematics **23**, European Mathematical Society (EMS), Zürich, 2015.
- [45] A. M. REYNOLDS AND C. J. RHODES, The Lévy flight paradigm: random search patterns and mechanisms, *Ecology* **90(4)** (2009), 877–887. DOI: 10.1890/08-0153.1.
- [46] X. ROS-OTON AND J. SERRA, The Dirichlet problem for the fractional Laplacian: regularity up to the boundary, *J. Math. Pures Appl. (9)* **101(3)** (2014), 275–302. DOI: 10.1016/j.matpur.2013.06.003.
- [47] X. ROS-OTON AND J. SERRA, Boundary regularity for fully nonlinear integro-differential equations, *Duke Math. J.* (to appear).
- [48] X. ROS-OTON AND J. SERRA, Regularity theory for general stable operators, Preprint (2014). [arXiv:1412.3892](#).
- [49] X. ROS-OTON, J. SERRA, AND E. VALDINOCI, Pohozaev identities for anisotropic integro-differential operators, Preprint (2015). [arXiv:1502.01431](#).
- [50] X. ROS-OTON AND E. VALDINOCI, The Dirichlet problem for nonlocal operators with singular kernels: convex and non-convex domains, Preprint (2015). [arXiv:1502.00782](#).
- [51] G. SAMORODNITSKY AND M. S. TAQQU, “*Stable non-Gaussian random processes*”, Stochastic models with infinite variance, Stochastic Modeling **1**, Chapman & Hall, New York, 1994.
- [52] W. SCHOUTENS, “*Lévy Processes in Finance: Pricing Financial Derivatives*”, Wiley Series in Probability and Statistics, John Wiley & Sons, New York, 2003. DOI: 10.1002/0470870230.
- [53] R. W. SCHWAB AND L. SILVESTRE, Regularity for parabolic integro-differential equations with very irregular kernels, Preprint (2014). [arXiv:1412.3790](#).
- [54] J. SERRA,  $C^{\sigma+\alpha}$  regularity for concave nonlocal fully nonlinear elliptic equations with rough kernels, Preprint (2014). [arXiv:1405.0930](#).
- [55] J. SERRA, Regularity for fully nonlinear nonlocal parabolic equations with rough kernels, *Calc. Var. Partial Differential Equations* **54(1)** (2015), 615–629. DOI: 10.1007/s00526-014-0798-6.

- [56] R. SERVADEI AND E. VALDINOCI, Weak and viscosity solutions of the fractional Laplace equation, *Publ. Mat.* **58(1)** (2014), 133–154. DOI: 10.5565/PUBLMAT-58114-06.
- [57] S. A. SILLING, Reformulation of elasticity theory for discontinuities and long-range forces, *J. Mech. Phys. Solids* **48(1)** (2000), 175–209. DOI: 10.1016/S0022-5096(99)00029-0.
- [58] E. M. STEIN, “*Singular integrals and differentiability properties of functions*”, Princeton Mathematical Series **30**, Princeton University Press, Princeton, N.J., 1970.
- [59] K. ZHOU AND Q. DU, Mathematical and numerical analysis of linear peridynamic models with nonlocal boundary conditions, *SIAM J. Numer. Anal.* **48(5)** (2010), 1759–1780. DOI: 10.1137/090781267.

The University of Texas at Austin  
Department of Mathematics  
2515 Speedway  
Austin, TX 78751  
USA  
*E-mail address:* `ros.oton@math.utexas.edu`

Primera versió rebuda el 16 d'abril de 2015,  
darrera versió rebuda el 27 d'agost de 2015.