# CORRIGENDUM TO "SHAPE AND PERIOD OF LIMIT CYCLES BIFURCATING FROM A CLASS OF HAMILTONIAN PERIOD ANNULUS"[NONLINEAR ANAL. 81 (2013) 130-148] 

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#### Abstract

In our paper [1] we are concerned with the problem of shape and period of isolated periodic solutions of perturbed analytic radial Hamiltonian vector fields in the plane. Actually, there is a mistake in the formula of the first order approximation of the period given in Corollary 4. Here we give its proper drafting.


Corollary 4 of [1] provides the expression for the period function of the limit cycles bifurcating from the period annulus of the class of radial Hamiltonian differential equations in the plane given by

$$
\left\{\begin{array}{l}
\dot{x}=-\frac{\partial}{\partial y} H(x, y)+\varepsilon P(x, y, \varepsilon),  \tag{1}\\
\dot{y}=\frac{y}{\partial x} H(x, y)+\varepsilon Q(x, y, \varepsilon),
\end{array}\right.
$$

where $H(x, y), P(x, y, \varepsilon)$ and $Q(x, y, \varepsilon)$ are analytic functions and $\varepsilon$ is a small parameter. We assume that this Hamiltonian vector field has a continuum of periodic orbits around the origin. In $(r, \theta)$-polar coordinates, $H$ only depends on $r, H=H(r)$ and the differential system (11) is written as the one-form

$$
\begin{equation*}
d H+\sum_{i=1}^{\infty} \varepsilon^{i}\left(S_{i}(r, \theta) d r-R_{i}(r, \theta) d \theta\right)=0 \tag{2}
\end{equation*}
$$

The general expression for the period of the isolated periodic orbits of (2) given in [1] is correct but there is a mistake in the first order term in its series in $\varepsilon$. The corrected expression is done in the corollary below. We remark that only the general expression for the period is used in [1]. Hence, all the expressions for the period described in the applications are correct.

Corollary 4. Let us assume the hypotheses of Theorem 1. Then, the period of the periodic solution, $r(\theta ; \rho, \varepsilon)$, of equation (2) given in Theorem 3 satisfies

$$
\begin{equation*}
T(\varepsilon ; \rho)=\int_{0}^{2 \pi} \frac{r(\theta ; \rho, \varepsilon)}{H^{\prime}(r(\theta ; \rho, \varepsilon))+\sum_{i=1}^{\infty} \varepsilon^{i} S_{i}(r(\theta ; \rho, \varepsilon), \theta)} d \theta . \tag{3}
\end{equation*}
$$

In particular,

$$
\begin{equation*}
T(\varepsilon ; \rho)=\frac{2 \pi \rho}{H^{\prime}(\rho)}+\varepsilon\left(\frac{2 \pi F_{2}(\rho)\left(\rho H^{\prime \prime}(\rho)-H^{\prime}(\rho)\right)}{\left(H^{\prime}(\rho)\right)^{2} F_{1}^{\prime}(\rho)}-\frac{\rho}{\left(H^{\prime}(\rho)\right)^{2}} \int_{0}^{2 \pi} S_{1}(\rho, \theta) d \theta\right)+O\left(\varepsilon^{2}\right) \tag{4}
\end{equation*}
$$

Proof. Assume that, in equation (1),

$$
P(x, y, \varepsilon)=\sum_{i=0}^{\infty} \varepsilon^{i} P_{i}(x, y), \quad Q(x, y, \varepsilon)=\sum_{i=0}^{\infty} \varepsilon^{i} Q_{i}(x, y),
$$

where $P_{i}, Q_{i}$ are analytic functions. Hence, in polar coordinates $(x, y)=(r \cos \theta, r \sin \theta)$, the angle variation of equation (1) writes as

$$
\begin{equation*}
\frac{d \theta}{d t}=\frac{1}{r} \frac{\partial H}{\partial r}(r, \theta)+\frac{1}{r} \sum_{i=1}^{\infty} \varepsilon^{i} S_{i}(r, \theta) \tag{5}
\end{equation*}
$$

where

$$
S_{i}(r, \theta)=\cos \theta Q_{i}(r \cos \theta, r \sin \theta)-\sin \theta P_{i}(r \cos \theta, r \sin \theta)
$$

Since equation (1) is of radial Hamiltonian type we write $H^{\prime}(r)=\frac{\partial H}{\partial r}(r, \theta)$ and, hence, (5) writes as the 1-form

$$
\begin{equation*}
d t=\frac{r d \theta}{H^{\prime}(r)+\sum_{i=1}^{\infty} \varepsilon^{i} S_{i}(r, \theta)} \tag{6}
\end{equation*}
$$

Expression (3) follows from (6) by direct integration on $t$ and by taking into account that $r=r(\theta ; \rho, \varepsilon)$.

To obtain formula (4), first we develop the integrand of expression (3), up to first order, in $\varepsilon$-power series and we get

$$
\begin{equation*}
T(\varepsilon ; \rho)=\frac{2 \pi \rho}{H_{r}^{\prime}(\rho)}-\frac{\varepsilon}{\left(H_{r}^{\prime}(\rho)\right)^{2}} \int_{0}^{2 \pi}\left(\left(\rho H^{\prime \prime}(\rho)-H^{\prime}(\rho)\right) r_{1}(\theta)+\rho S_{1}(\rho, \theta)\right) d \theta \tag{7}
\end{equation*}
$$

where $r_{1}(\theta)$ is given by the development of $r(\theta ; \rho, \varepsilon)$ in its $\varepsilon$-power series

$$
r(\theta ; \rho, \varepsilon)=r_{0}(\theta)+\varepsilon r_{1}(\theta)+\varepsilon^{2} r_{2}(\theta)+\varepsilon^{3} r_{3}(\theta)+\cdots .
$$

Finally, by using the expression of $r_{1}(\theta)$, given in Theorem 1, and Theorem 3.(i) of [1], formula (4) follows.

## References

[1] Prohens, R. and Torregrosa, J. Shape and period of limit cycles bifurcating from a class of Hamiltonian period annulus. Nonlinear Anal. Ser. A, 81 (2013) 1307-148.
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