

OSCILLATORY PHENOMENA FOR HIGHER-ORDER FRACTIONAL LAPLACIANS

NICOLA ABATANGELO AND SVEN JAROHS

Abstract: We collect some peculiarities of higher-order fractional Laplacians $(-\Delta)^s$, $s > 1$, with special attention to the range $s \in (1, 2)$, which show their oscillatory nature. These include the failure of the polarization and Pólya–Szegő inequalities and the explicit example of a domain with sign-changing first eigenfunction. In spite of these fluctuating behaviours, we prove how the Faber–Krahn inequality still holds for any $s > 1$ in dimension one.

2020 Mathematics Subject Classification: Primary: 35A23, 35G15, 35P05; Secondary: 47A75, 49R20, 49R50.

Key words: polarization inequality, Pólya–Szegő inequality, first eigenfunction, positivity-preserving properties, Faber–Krahn inequality.

1. Introduction

The fractional Laplacian operator, usually denoted by $(-\Delta)^s$, $s \in (0, 1)$, is the nonlocal integral operator

$$(1) \quad (-\Delta)^s u(x) = c_{n,s} \lim_{\varepsilon \downarrow 0} \int_{\{y \in \mathbb{R}^n : |y| > \varepsilon\}} \frac{u(x) - u(x+y)}{|y|^{n+2s}} dy.$$

It is naturally associated to the stochastic analysis of some Lévy processes, in particular of the α -stable ones, since it arises as their infinitesimal generator. For this reason, it has been used in models from mathematical finance, population dynamics, quantum mechanics, and electrostatics, among others. We refer to [6, 10, 11, 17] for an introduction to the basic features and a more comprehensive list of applications of the operator. From an abstract perspective, the family $\{(-\Delta)^s\}_{s \in (0,1)}$ can be thought of as a collection of operators interpolating the identity and the Laplacian; indeed, at least for $u \in C_c^\infty(\mathbb{R}^n)$,

$$\lim_{s \downarrow 0} (-\Delta)^s u = u,$$

$$\lim_{s \uparrow 1} (-\Delta)^s u = -\Delta u;$$

see [11, Proposition 4.4].

Here, we are mainly concerned with powers of the Laplace operator greater than 1, which we call *higher-order fractional Laplacians*. Although the integral expression in (1) cannot be extended to $s > 1$ in a straightforward way, it is possible to define $(-\Delta)^s$, $s > 1$, as the pseudo-differential operator

$$(2) \quad \mathcal{F}[(-\Delta)^s u](\xi) = |\xi|^{2s} \mathcal{F}u(\xi) \quad \text{for } \xi \in \mathbb{R}^n, u \in C_c^\infty(\mathbb{R}^n),$$

where \mathcal{F} denotes the Fourier transform. For $s \in (0, 1)$, the above and (1) coincide by properly choosing the normalizing constant $c_{n,s}$; see [31, Chapter 5, Lemma 25.3, and Theorem 26.1]. With definition (2), the family $\{(-\Delta)^s\}_{s > 0}$ interpolates the (local) *polyharmonic operators* – or also *poly-Laplacians* – $(-\Delta)^s$, $s \in \mathbb{N}$, for which we refer to the monograph [18].

Boundary problems driven by a poly-Laplacian require several boundary conditions in order to be well posed, and this is due to the high order of the operator. One way of prescribing boundary conditions is to impose that the solution u satisfies

$$u = \frac{\partial u}{\partial \nu} = \dots = \frac{\partial^{s-1} u}{\partial \nu^{s-1}} = 0 \quad \text{at the boundary, for } s \in \mathbb{N},$$

where ν denotes the normal unit vector to the boundary. These go under the name *Dirichlet conditions* and they represent the pointwise analogue of variationally solving equations in the Sobolev space H_0^s .

When coupled with Dirichlet conditions, poly-Laplacians present an oscillatory behaviour. This may be exemplified by the failure of the weak maximum principle

$$(3) \quad (-\Delta)^s u \geq 0 \quad \text{in } \Omega \quad (u \in H_0^s(\Omega), s \in \mathbb{N} \setminus \{1\}) \not\Rightarrow u \geq 0 \quad \text{in } \Omega.$$

It was first noted in 1908 [21] that, in dimension $n = 2$, there are annular domains for which (3) occurs for the bi-Laplacian $s = 2$. It was then conjectured that (3) was *not* the case for convex domains: this is known in the literature as the *Boggio–Hadamard conjecture*. But this conjecture was proved to be false [13] for some rectangular domains and, later on, several other counterexamples were built; these include for $s, n = 2$: an infinite strip [14], (most) infinite wedges [32], the punctured disk [28], and eccentric ellipses [16]. As to this last counterexample, it is worth mentioning that, on top of the simple geometry of the domain, it is possible to provide completely elementary counterexamples in terms of polynomials [33]: this highlights how the above mentioned oscillations are deeply written in the nature of higher-order operators and they do not arise as singular phenomena. For $s = 3$ we refer to [34] and for $s = 4$ to [35] for elementary explicit counterexamples. For $s \in \mathbb{N}$ even and any space dimension, the lack of maximum principles can be deduced by the analysis of the oscillations of the first eigenfunction carried out in [24].

Another interesting (and extremely difficult) problem for the bi-Laplacian with Dirichlet conditions is the *Faber–Krahn inequality*. It states that the first eigenvalue of any domain is larger than the first eigenvalue of the ball with the same measure. This has been proved independently for $n = 2$ in [27] and for $n = 2, 3$ in [7], while it remains completely open for higher dimensions. Both [27] and [7] are based on the previous analysis carried out in [37] which, in turn, borrows the idea of *Talenti’s principle* [36, Theorem 1].

As the collection $\{(-\Delta)^s\}_{s \in (1,2)}$ connects the Laplacian with the bi-Laplacian, one could legitimately wonder in which way this oscillatory behaviour appears in the transition from $s = 1$ to $s = 2$. Our analysis here aims at contributing to answer this question, by showing that as soon as $s > 1$ oscillatory phenomena emerge to a large extent.

First of all, let us recall that, on the fractional and nonlocal side $s \in (1, \infty) \setminus \mathbb{N}$, it was first proved in [3] that “almost” no disconnected domain can satisfy the weak maximum principle whenever $s \in (2k - 1, 2k)$ with $k \in \mathbb{N}$: the prototypical domain here can be thought of as the union of two disjoint balls (bear in mind that the disconnection of the domain is overruled by the nonlocality of the operator). A somewhat more precise analysis in [2, Theorem 1.10] studied the sign of the Green function of such two-ball domain, noticing in particular how the Green function is positive (if the balls have equal radii) for $s \in (2k, 2k + 1)$, $k \in \mathbb{N}$, and thus the above mentioned counterexample cannot extend also to this range of s . Later, counterexamples on eccentric ellipsoids for $s \in (1, \sqrt{3} + 3/2)$ were constructed in [5], therefore covering the exponents $s \in (2, 3)$.

Another interesting property found in [26] deals with the Dirichlet energy¹

$$(4) \quad \mathcal{E}_s(u, u) = \int_{\mathbb{R}^n} |\xi|^{2s} |\mathcal{F}u(\xi)|^2 d\xi$$

of the Nemytskii operator $u \mapsto |u|$: it holds that

$$\mathcal{E}_s(u, u) \leq \mathcal{E}_s(|u|, |u|) \quad \text{for } s \in (1, 3/2),$$

where the equality holds only if u is of constant sign. This is in particular equivalent to

$$\mathcal{E}_s(u_+, u_-) \geq 0 \quad \text{for } s \in (1, 3/2),$$

where u_+, u_- denote respectively the positive and the negative parts of u .

The purpose of the present note is to collect a series of remarks on the behaviour of higher-order fractional Laplacians. In particular, we show that:

- The polarization inequality for $s \in (1, 3/2)$ is reversed with respect to what happens for $s \in (0, 1]$ – see Definition 3.1 and Theorem 3.8. The polarization of a function is a rearrangement with respect to a hyperplane and it is a useful tool in proving symmetry of solutions to elliptic and parabolic problems: we refer to [41] and the many references therein for a survey on this matter. Among the applications of the polarization, it is worth mentioning that it can approximate several other notions of rearrangement (including the spherical decreasing one) and that, as a consequence, a number of rearrangement inequalities can be deduced from the polarization inequality; see [39].
- The Pólya–Szegő inequality fails for $s \in (1, \infty)$; see Subsection 4.2. Contrarily to what happens for the polarization, the Pólya–Szegő inequality is not reversed; see Subsection 4.1. This inequality compares the energy of the spherical decreasing rearrangement (or, also, Schwarz symmetrization) – see Definition 4.1 – of a function with the energy of the original function. For $s \in (0, 1)$ the Pólya–Szegő inequality has been proved in [29, equation (14)]; recall that this inequality implies the Faber–Krahn inequality.
- On a two-ball domain, the first eigenvalue is simple and the first eigenfunction is symmetric and positive for $s \in (2k, 2k + 1)$, $k \in \mathbb{N}$, whilst it is antisymmetric and positive on one ball (and negative on the other) for $s \in (2k - 1, 2k)$, $k \in \mathbb{N}$; see Theorem 5.4.
- Again on a two-ball domain, when the balls have equal radii, maximum principles are recovered for antisymmetric “positive” data for $s \in (2k, 2k + 1)$, $k \in \mathbb{N}$, and for symmetric positive data for $s \in (2k - 1, 2k)$, $k \in \mathbb{N}$; see Proposition 5.5.
- Despite the results concerning the polarization and the Pólya–Szegő inequalities, the Faber–Krahn inequality still holds for any $s \in (1, \infty)$ at least in dimension $n = 1$; see Theorems 6.4 and 6.6.

2. Notations and reminders about higher-order fractional Laplacians

In the following we will need some notations and known facts about the analysis of higher-order fractional Laplacians, which we recall here below for future reference.

We will denote by $[\cdot]$ the integer part of a real number, *i.e.*,

$$[s] = \max\{m \in \mathbb{Z} : m < s\}, \quad s \in \mathbb{R}.$$

¹This could be equivalently introduced as the quadratic form which is naturally associated to $(-\Delta)^s$ as defined in (2). For $s = 0$ it coincides with the square of the L^2 norm.

Note that, with this definition, $\lfloor s \rfloor = s - 1$ whenever $s \in \mathbb{N}$. If \mathbb{N} denotes the set of positive integer numbers, then $\mathbb{N}_0 = \mathbb{N} \cup \{0\}$ and $2\mathbb{N}_0, 2\mathbb{N}_0 + 1$ denote respectively the sets of even and odd nonnegative integers. The symbols \vee and \wedge will denote respectively the max and min operations between real numbers, namely

$$a \vee b = \max\{a, b\}, \quad a \wedge b = \min\{a, b\}, \quad \text{for } a, b \in \mathbb{R}.$$

By ω_n we mean the volume of the unit ball in \mathbb{R}^n ,

$$\omega_n = \frac{2\pi^{n/2}}{n\Gamma(n/2)}.$$

Given a measurable $A \subset \mathbb{R}^n$, we denote by $\mathbf{1}_A$ its characteristic function.

The fractional Laplacian $(-\Delta)^s$ as defined in (2) admits the pointwise representation (see [4])

$$(-\Delta)^s u(x) := \frac{\kappa_{n,s}}{2} \int_{\mathbb{R}^n} \frac{\delta_{\lfloor s \rfloor + 1} u(x, y)}{|y|^{n+2s}} dy,$$

where

$$\delta_{\lfloor s \rfloor + 1} u(x, y) := \sum_{k=-\lfloor s \rfloor - 1}^{\lfloor s \rfloor + 1} (-1)^k \binom{2\lfloor s \rfloor + 2}{\lfloor s \rfloor + 1 - k} u(x + ky) \quad \text{for } x, y \in \mathbb{R}^n$$

is a finite difference of order $2\lfloor s \rfloor + 2$, and $\kappa_{n,s}$ is the positive constant given by

$$\kappa_{n,s} := \begin{cases} \frac{2^{2s}\Gamma(n/2 + s)}{\pi^{n/2}\Gamma(-s)} \left(\sum_{k=1}^{\lfloor s \rfloor + 1} (-1)^k \binom{2\lfloor s \rfloor + 2}{\lfloor s \rfloor + 1 - k} k^{2s} \right)^{-1}, & s \notin \mathbb{N}, \\ \frac{2^{2s}\Gamma(n/2 + s) s!}{2\pi^{n/2}} \left(\sum_{k=2}^{s+1} (-1)^{k-s+1} \binom{2s+2}{s+1-k} k^{2s} \ln(k) \right)^{-1}, & s \in \mathbb{N}. \end{cases}$$

Recalling the quadratic form in (4), the fractional Sobolev space $H^s(\mathbb{R}^n)$ is defined as

$$H^s(\mathbb{R}^n) = \{u \in L^2(\mathbb{R}^n) : \mathcal{E}_s(u, u) < \infty\}.$$

On a bounded open domain $\Omega \subset \mathbb{R}^n$ the Sobolev space associated with homogeneous boundary conditions is

$$\mathcal{H}_0^s(\Omega) = \{u \in H^s(\mathbb{R}^n) : u = 0 \text{ in } \mathbb{R}^n \setminus \overline{\Omega}\}.$$

The bilinear form induced by (4) is

$$\mathcal{E}_s(u, v) = \int_{\mathbb{R}^n} |\xi|^{2s} \mathcal{F}u(\xi) \overline{\mathcal{F}v(\xi)} d\xi$$

and it admits also the following equivalent representations:

$$(5) \quad \mathcal{E}_s(u, v) = \begin{cases} \mathcal{E}_{s-\lfloor s \rfloor}((-\Delta)^{\lfloor s \rfloor/2} u, (-\Delta)^{\lfloor s \rfloor/2} v) & \text{if } \lfloor s \rfloor \in 2\mathbb{N}_0, \\ \mathcal{E}_{s-\lfloor s \rfloor}(\nabla(-\Delta)^{(\lfloor s \rfloor - 1)/2} u, \nabla(-\Delta)^{(\lfloor s \rfloor - 1)/2} v) & \text{if } \lfloor s \rfloor \in 2\mathbb{N}_0 + 1. \end{cases}$$

A consequence of this representation is that, if $u, v \in H^s(\mathbb{R}^n)$ with $uv \equiv 0$ in \mathbb{R}^n , then

$$(6) \quad \mathcal{E}_s(u, v) = (-1)^{\lfloor s \rfloor + 1} \frac{c_{n,s}}{2} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \frac{u(x)v(y)}{|x-y|^{n+2s}} dx dy,$$

where

$$c_{n,s} = \frac{2^{2s}\Gamma(n/2 + s)}{\pi^{n/2}|\Gamma(-s)|} > 0;$$

see [3, Lemma 4.4] and [26, Theorem 2]. Equation (6) should be borne in mind, as it will play a crucial role in the following.

3. A (reversed) polarization inequality

In the following, let Σ be an open half-space in \mathbb{R}^n . Denote by $\tau_\Sigma: \mathbb{R}^n \rightarrow \mathbb{R}^n$ the reflection around $\partial\Sigma$.

Definition 3.1. The *polarization* of a function $u: \mathbb{R}^n \rightarrow \mathbb{R}$ is given by

$$u_\Sigma: \mathbb{R}^n \longrightarrow \mathbb{R}, \quad u_\Sigma(x) = \begin{cases} u(x) \vee u(\tau_\Sigma(x)) & \text{for } x \in \Sigma, \\ u(x) \wedge u(\tau_\Sigma(x)) & \text{for } x \in \mathbb{R}^n \setminus \Sigma. \end{cases}$$

Note that we may also write

$$(7) \quad u_\Sigma(x) = \frac{u(x) + u(\tau_\Sigma(x))}{2} \begin{cases} + \frac{|u(x) - u(\tau_\Sigma(x))|}{2} & \text{for } x \in \Sigma, \\ - \frac{|u(x) - u(\tau_\Sigma(x))|}{2} & \text{for } x \in \mathbb{R}^n \setminus \Sigma, \end{cases}$$

and it holds that

$$(8) \quad u_\Sigma + u_\Sigma \circ \tau_\Sigma = u + u \circ \tau_\Sigma \quad \text{in } \mathbb{R}^n.$$

Recall that $\mathcal{E}_s(u, u) = \mathcal{E}_s(u_\Sigma, u_\Sigma)$ for $s = 0, 1$ (see, e.g., [41]) whereas, for $s \in (0, 1)$, we have

$$(9) \quad \mathcal{E}_s(u_\Sigma, u_\Sigma) \leq \mathcal{E}_s(u, u),$$

where the inequality is strict if $u_\Sigma \neq u$ (see, e.g., [8, 19, 40]). In the following, we consider $s \in (1, 3/2)$ and show that (9) is reversed; see Theorem 3.8 below. The fact that $u_\Sigma \in H^s(\mathbb{R}^n)$ whenever $u \in H^s(\mathbb{R}^n)$ follows immediately from (7) and [9, Théorème 1].

Lemma 3.2. *Let $s \in (0, 1)$, Σ an open half-space in \mathbb{R}^n , and $v \in H^s(\mathbb{R}^n)$ such that $v \circ \tau_\Sigma = -v$, i.e., v is antisymmetric with respect to $\partial\Sigma$. Then $v\mathbf{1}_\Sigma \in \mathcal{H}_0^s(\Sigma)$ and*

$$\mathcal{E}_s(v\mathbf{1}_\Sigma, v\mathbf{1}_\Sigma) \leq \mathcal{E}_s(v, v).$$

Proof: The statement follows from the proof of [23, Lemma 3.2] (see also [22, Lemma 3.2]). □

Proposition 3.3. *Let $s \in (1, 3/2)$, Σ an open half-space in \mathbb{R}^n , and $v \in H^s(\mathbb{R}^n)$ such that $v \circ \tau_\Sigma = -v$, i.e., v is antisymmetric w.r.t. $\partial\Sigma$. Then $v\mathbf{1}_\Sigma \in \mathcal{H}_0^s(\Sigma)$.*

Proof: By translation and rotation invariance, we may assume that $\Sigma = \{x_1 > 0\}$, so that $\tau_\Sigma(x) = \tau_\Sigma(x_1, x') = (-x_1, x')$. Note that $v\mathbf{1}_\Sigma \in \mathcal{H}_0^1(\Sigma)$. Next recall that by [20, Corollary 1.4.4.5] we have $\mathcal{H}_0^{s-1}(\Sigma) = H^{s-1}(\Sigma)$ for $s \in (1, 3/2)$. In particular, for any $i = 1, \dots, n$, we have $(\partial_i v)\mathbf{1}_\Sigma \in H^{s-1}(\Sigma)$ and thus also $\partial_i(v\mathbf{1}_\Sigma) = (\partial_i v)\mathbf{1}_\Sigma \in \mathcal{H}_0^{s-1}(\Sigma)$. □

Remark 3.4. We note that the conclusion of Proposition 3.3 does not hold for $s \geq 3/2$. Indeed, for example in the case $n = 1$ and $s = 2$, if $\Sigma = \{x_1 > 0\}$ and $v \in H^2(\mathbb{R})$ is odd, then in general $v'(0) \neq 0$ and thus $v\mathbf{1}_\Sigma$ cannot be an element of $\mathcal{H}_0^2(\Sigma)$. More generally, one can check [25, Example 3.1(i)] or [26, Example 1].

Lemma 3.5. *Let $s \in (1, 2)$, Σ an open half-space in \mathbb{R}^n , and $w_1, w_2 \in \mathcal{H}_0^s(\Sigma)$ with $w_1, w_2 \geq 0$ and $w_1 w_2 \equiv 0$ in \mathbb{R}^n . Then*

$$\mathcal{E}_s(w_1, w_2 \circ \tau_\Sigma) \leq \mathcal{E}_s(w_1, w_2)$$

and the inequality is strict if w_1 and w_2 are nontrivial.

Proof: By density, we may think of $w_1, w_2 \in C_c^\infty(\Sigma)$ with disjoint supports. Moreover, since $w_1 w_2 \equiv 0$ in \mathbb{R}^n and $w_i \mathbf{1}_{\mathbb{R}^n \setminus \Sigma} \equiv 0$ in \mathbb{R}^n , $i = 1, 2$, we also have $w_1(w_2 \circ \tau_\Sigma) \equiv 0$ in \mathbb{R}^n . Hence, (6) implies

$$\begin{aligned} \mathcal{E}_s(w_1, w_2 \circ \tau_\Sigma) &= \frac{c_{n,s}}{2} \int_\Sigma \int_{\mathbb{R}^n \setminus \Sigma} \frac{w_1(x) w_2(\tau_\Sigma(y))}{|x - y|^{n+2s}} dy dx \\ &= \frac{c_{n,s}}{2} \int_\Sigma \int_\Sigma \frac{w_1(x) w_2(y)}{|x - \tau_\Sigma(y)|^{n+2s}} dy dx \\ &\leq \frac{c_{n,s}}{2} \int_\Sigma \int_\Sigma \frac{w_1(x) w_2(y)}{|x - y|^{n+2s}} dy dx = \mathcal{E}_s(w_1, w_2), \end{aligned}$$

since $|x - \tau_\Sigma(y)| \geq |x - y|$ for $x, y \in \Sigma$. □

Lemma 3.6. *Let $s \in (1, 3/2)$, Σ an open half-space in \mathbb{R}^n , and $w \in H^s(\mathbb{R}^n)$ such that $w \circ \tau_\Sigma = -w$, i.e., w is antisymmetric with respect to $\partial\Sigma$. Then*

$$(10) \quad \mathcal{E}_s(w, w) \leq \mathcal{E}_s(|w| \mathbf{1}_\Sigma - |w| \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}, |w| \mathbf{1}_H - |w| \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}).$$

Moreover, the inequality is strict if $w \mathbf{1}_\Sigma$ is sign-changing.

Proof: By Proposition 3.3 we have $w \mathbf{1}_\Sigma \in \mathcal{H}_0^s(\Sigma)$ and hence also $(w \mathbf{1}_\Sigma)^\pm = w^\pm \mathbf{1}_\Sigma \in \mathcal{H}_0^s(\Sigma)$ by [9, Théorème 1], since $s < 3/2$. Splitting $w = w^+ \mathbf{1}_\Sigma - w^- \mathbf{1}_\Sigma + w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma} - w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}$ and $|w| = w^+ + w^-$ we have that (10) is equivalent to

$$\begin{aligned} & -\mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^- \mathbf{1}_\Sigma) + \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) - \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) \\ & -\mathcal{E}_s(w^- \mathbf{1}_\Sigma, w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) + \mathcal{E}_s(w^- \mathbf{1}_\Sigma, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) - \mathcal{E}_s(w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) \\ & \leq \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^- \mathbf{1}_\Sigma) - \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) - \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) \\ & \quad - \mathcal{E}_s(w^- \mathbf{1}_\Sigma, w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) - \mathcal{E}_s(w^- \mathbf{1}_\Sigma, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) + \mathcal{E}_s(w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}), \end{aligned}$$

or, simply,

$$\begin{aligned} \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) + \mathcal{E}_s(w^- \mathbf{1}_\Sigma, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}) \\ \leq \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^+ \mathbf{1}_\Sigma) + \mathcal{E}_s(w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}, w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma}). \end{aligned}$$

Using the antisymmetry of w , one has that

$$w^+ \mathbf{1}_{\mathbb{R}^n \setminus \Sigma} = (w^- \mathbf{1}_\Sigma) \circ \tau_\Sigma \quad \text{and} \quad w^- \mathbf{1}_{\mathbb{R}^n \setminus \Sigma} = (w^+ \mathbf{1}_\Sigma) \circ \tau_\Sigma,$$

so that we only need to show

$$\mathcal{E}_s(w^+ \mathbf{1}_\Sigma, (w^- \mathbf{1}_\Sigma) \circ \tau_\Sigma) \leq \mathcal{E}_s(w^+ \mathbf{1}_\Sigma, w^- \mathbf{1}_\Sigma),$$

since the scalar product is invariant under reflections. This last inequality holds in view of Lemma 3.5. □

Proposition 3.7. *Let $s \in (1, 3/2)$, Σ an open half-space in \mathbb{R}^n , and $u \in H^s(\mathbb{R}^n)$. Then*

$$\mathcal{E}_s(u_\Sigma, u_\Sigma \circ \tau_\Sigma) \leq \mathcal{E}_s(u, u \circ \tau_\Sigma).$$

Moreover, this inequality is strict if $u_\Sigma \neq u$.

Proof: First note that

$$4\mathcal{E}_s(u, u \circ \tau_\Sigma) = \mathcal{E}_s(u + u \circ \tau_\Sigma, u + u \circ \tau_\Sigma) - \mathcal{E}_s(u - u \circ \tau_\Sigma, u - u \circ \tau_\Sigma).$$

Using this observation and (7), it follows that

$4\mathcal{E}_s(u_\Sigma, u_\Sigma \circ \tau_\Sigma) - 4\mathcal{E}_s(u, u \circ \tau_\Sigma) = \mathcal{E}_s(u - u \circ \tau_\Sigma, u - u \circ \tau_\Sigma) - \mathcal{E}_s(u_\Sigma - u_\Sigma \circ \tau_\Sigma, u_\Sigma - u_\Sigma \circ \tau_\Sigma)$
 since $u + u \circ \tau_\Sigma = u_\Sigma + u_\Sigma \circ \tau_\Sigma$ by definition. Noting that $u - u \circ \tau_\Sigma$ is antisymmetric with respect to $\partial\Sigma$ and that

$$u_\Sigma - u_\Sigma \circ \tau_\Sigma = |u - u \circ \tau_\Sigma| \mathbf{1}_\Sigma - |u - u \circ \tau_\Sigma| \mathbf{1}_{\mathbb{R}^n \setminus \Sigma},$$

the assertion follows from Lemma 3.6. □

Theorem 3.8. *Let $s \in (1, 3/2)$ and Σ an open half-space in \mathbb{R}^n . Then for all $u \in H^s(\mathbb{R}^n)$ we have $u_\Sigma \in H^s(\mathbb{R}^n)$ and*

$$\mathcal{E}_s(u_\Sigma, u_\Sigma) \geq \mathcal{E}_s(u, u).$$

Moreover, this inequality is strict if $u_\Sigma \neq u$.

Proof: Recall identity (8). Note that

$$\begin{aligned} \mathcal{E}_s(u_\Sigma, u_\Sigma) &= \frac{1}{2} \mathcal{E}_s(u_\Sigma, u_\Sigma) + \frac{1}{2} \mathcal{E}_s(u_\Sigma \circ \tau_\Sigma, u_\Sigma \circ \tau_\Sigma) \\ &= \frac{1}{2} \mathcal{E}_s(u_\Sigma + u_\Sigma \circ \tau_\Sigma, u_\Sigma + u_\Sigma \circ \tau_\Sigma) - \mathcal{E}_s(u_\Sigma, u_\Sigma \circ \tau_\Sigma) \\ &= \frac{1}{2} \mathcal{E}_s(u + u \circ \tau_\Sigma, u + u \circ \tau_\Sigma) - \mathcal{E}_s(u_\Sigma, u_\Sigma \circ \tau_\Sigma) \\ &= \mathcal{E}_s(u, u) + \mathcal{E}_s(u, u \circ \tau_\Sigma) - \mathcal{E}_s(u_\Sigma, u_\Sigma \circ \tau_\Sigma). \end{aligned}$$

Note that we have used (8). The claim hence follows from Proposition 3.7. □

4. The Pólya–Szegő inequality

We start by recalling the definition of the symmetric decreasing rearrangement. We retrieve the definitions in [36].

Definition 4.1. Given a measurable function $u: \mathbb{R}^n \rightarrow \mathbb{R}$ the *decreasing rearrangement* of u is the function

$$\begin{aligned} \tilde{u}: [0, \infty) &\longrightarrow [0, \infty) \\ t &\longmapsto \inf\{\tau \geq 0 : |\{ |u| > \tau \}| < t\} \end{aligned}$$

and the *spherical decreasing rearrangement* is

$$\begin{aligned} u^*: \mathbb{R}^n &\longrightarrow [0, \infty) \\ x &\longmapsto \tilde{u}(\omega_n |x|^n). \end{aligned}$$

4.1. An example. Consider $v \in \mathcal{H}_0^s(B_1)$, nonnegative, radial, and radially decreasing. Note that we can also write $v(x) = w(|x|)$, $x \in \mathbb{R}^n$, for some $w: [0, \infty) \rightarrow [0, \infty)$.

Fix $x_0 \in \mathbb{R}^n$, $|x_0| > 2$, and define

$$u(x) := v(x) + v_{a,x_0}(x), \quad v_{a,x_0}(x) := v(ax - x_0), \quad x \in \mathbb{R}^n, a \geq 1.$$

For $t \in \mathbb{R}$, consider the superlevels of u , i.e., $A_t = \{x \in \mathbb{R}^n : u(x) > t\}$. It holds that

$$|A_t| = |\{x \in \mathbb{R}^n : u(x) > t\}| = |\{x \in \mathbb{R}^n : v(x) > t\}| + |\{x \in \mathbb{R}^n : v(ax) > t\}|,$$

so that the symmetrization A_t^* of A_t coincides with

$$A_t^* = B(0, \omega_n^{-1/n} (|\{x \in \mathbb{R}^n : v(x) > t\}| + |\{x \in \mathbb{R}^n : v(ax) > t\}|)^{1/n}).$$

We now compute u^* , the spherical decreasing rearrangement of u . By definition,

$$u^*(y) = \inf\{t \in \mathbb{R} : y \in A_t^*\},$$

where

$$\begin{aligned} y \in A_t^* & \text{ if and only if } |y|^n < \omega_n^{-1}|\{x \in \mathbb{R}^n : v(x) > t\}| + \omega_n^{-1}|\{x \in \mathbb{R}^n : v(ax) > t\}| \\ & = |\{z > 0 : w(z) > t\}|^n + |\{z > 0 : w(az) > t\}|^n \\ & = |w^{-1}(t)|^n + \left| \frac{w^{-1}(t)}{a} \right|^n = \left(1 + \frac{1}{a^n} \right) |w^{-1}(t)|^n, \end{aligned}$$

which implies

$$u^*(y) = w\left(\frac{a|y|}{(1+a^n)^{1/n}}\right) = v_{b,0}(y), \quad b := \frac{a}{(1+a^n)^{1/n}}, \quad \text{for } y \in \mathbb{R}^n.$$

Let us now compare the Sobolev energies:

$$\begin{aligned} \mathcal{E}_s(u, u) &= \mathcal{E}_s(v, v) + 2\mathcal{E}_s(v, v_{a,x_0}) + \mathcal{E}_s(v_{a,x_0}, v_{a,x_0}) \\ &= (1 + a^{2s-n})\mathcal{E}_s(v, v) + 2\mathcal{E}_s(v, v_{a,x_0}) \geq (1 + a^{2s-n})\mathcal{E}_s(v, v), \end{aligned}$$

$$\mathcal{E}_s(u^*, u^*) = \mathcal{E}_s(v_{b,0}, v_{b,0}) = b^{2s-n}\mathcal{E}_s(v, v) = \frac{a^{2s-n}}{(1+a^n)^{(2s-n)/n}}\mathcal{E}_s(v, v);$$

here we must note that

$$(11) \quad \frac{a^{2s-n}}{(1+a^n)^{(2s-n)/n}} < 1 + a^{2s-n}, \quad a > 1,$$

so that we find

$$(12) \quad \mathcal{E}_s(u^*, u^*) < \mathcal{E}_s(u, u).$$

Indeed, (11) is equivalent to

$$(1+a^n)^{1-2s/n} < 1 + (a^n)^{1-2s/n}, \quad a > 1,$$

which holds by the subadditivity of $t \mapsto t^\gamma$ for $t \geq 1, \gamma \leq 1$.

4.2. A counterexample. Let us consider $n = 1$, restrict $s \in (1, 3/2)$, and define for $M > 1$

$$u(x) := \begin{cases} 0 & \text{for } x \leq -2, \\ x+2 & \text{for } -2 < x \leq 0, \\ -x+2 & \text{for } 0 < x \leq 1, \\ 1 & \text{for } 1 < x \leq 2M-1, \\ -x+2M & \text{for } 2M-1 < x \leq 2M, \\ 0 & \text{for } x > 2M. \end{cases}$$

The spherical rearrangement of u is

$$u^*(x) := \begin{cases} -|x|+2 & \text{for } 0 \leq |x| \leq 1, \\ 1 & \text{for } 1 < |x| \leq M, \\ -|x|+M+1 & \text{for } M < |x| \leq M+1, \\ 0 & \text{for } |x| > M+1. \end{cases}$$

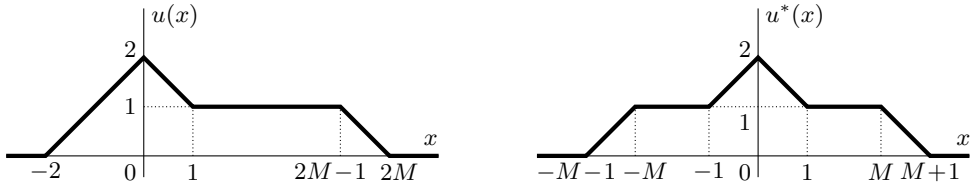


FIGURE 1. The graphs of u and u^* .

Let us now explicitly compute the Dirichlet energies associated to both functions. As to u , we have (we are going to extensively use translation invariance and (5))

$$\begin{aligned} \mathcal{E}_s(u, u) &= \mathcal{E}_{s-1}(u', u') \\ &= \mathcal{E}_{s-1}(\mathbf{1}_{(-2,0)} - \mathbf{1}_{(0,1)} - \mathbf{1}_{(2M-1,2M)}, \mathbf{1}_{(-2,0)} - \mathbf{1}_{(0,1)} - \mathbf{1}_{(2M-1,2M)}) \\ &= \mathcal{E}_{s-1}(\mathbf{1}_{(-2,-1)} + \mathbf{1}_{(-1,0)} - \mathbf{1}_{(0,1)} - \mathbf{1}_{(2M-1,2M)}, \\ &\quad \mathbf{1}_{(-2,-1)} + \mathbf{1}_{(-1,0)} - \mathbf{1}_{(0,1)} - \mathbf{1}_{(2M-1,2M)}) \\ &= 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(0,1)}) - 2\mathcal{E}_{s-1}(\mathbf{1}_{(-2,-1)}, \mathbf{1}_{(0,1)}) - 2\mathcal{E}_s(\mathbf{1}_{(-2,-1)}, \mathbf{1}_{(2M-1,2M)}) \\ &\quad - 2\mathcal{E}_{s-1}(\mathbf{1}_{(-1,0)}, \mathbf{1}_{(2M-1,2M)}) + 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M-1,2M)}), \end{aligned}$$

whereas

$$\begin{aligned} \mathcal{E}_s(u^*, u^*) &= \mathcal{E}_{s-1}(\mathbf{1}_{(-M-1,-M)} + \mathbf{1}_{(-1,0)} - \mathbf{1}_{(0,1)} - \mathbf{1}_{(M,M+1)}, \\ &\quad \mathbf{1}_{(-M-1,-M)} + \mathbf{1}_{(-1,0)} - \mathbf{1}_{(0,1)} - \mathbf{1}_{(M,M+1)}) \\ &= 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(0,1)}) + 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(M,M+1)}) - 4\mathcal{E}_{s-1}(\mathbf{1}_{(-1,0)}, \mathbf{1}_{(M,M+1)}) \\ &\quad - 2\mathcal{E}_{s-1}(\mathbf{1}_{(-M-1,-M)}, \mathbf{1}_{(M,M+1)}) - 2\mathcal{E}_{s-1}(\mathbf{1}_{(-1,0)}, \mathbf{1}_{(0,1)}). \end{aligned}$$

Summing up,

$$\begin{aligned} \mathcal{E}_s(u, u) - \mathcal{E}_s(u^*, u^*) &= -2\mathcal{E}_{s-1}(\mathbf{1}_{(-2,-1)}, \mathbf{1}_{(0,1)}) - 2\mathcal{E}_s(\mathbf{1}_{(-2,-1)}, \mathbf{1}_{(2M-1,2M)}) \\ &\quad - 2\mathcal{E}_{s-1}(\mathbf{1}_{(-1,0)}, \mathbf{1}_{(2M-1,2M)}) + 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M-1,2M)}) \\ &\quad - 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(M,M+1)}) + 4\mathcal{E}_{s-1}(\mathbf{1}_{(-1,0)}, \mathbf{1}_{(M,M+1)}) \\ &\quad + 2\mathcal{E}_{s-1}(\mathbf{1}_{(-M-1,-M)}, \mathbf{1}_{(M,M+1)}) + 2\mathcal{E}_{s-1}(\mathbf{1}_{(-1,0)}, \mathbf{1}_{(0,1)}) \\ &= -2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(-2,-1)}) - 2\mathcal{E}_s(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M+1,2M+2)}) \\ &\quad - 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M,2M+1)}) + 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M-1,2M)}) \\ &\quad - 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(M,M+1)}) + 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(M+1,M+2)}) \\ &\quad + 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M+1,2M+2)}) + 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(-1,0)}) \\ &= -2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(-2,-1)}) + 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(-1,0)}) \\ &\quad - 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M,2M+1)}) + 2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(2M-1,2M)}) \\ &\quad - 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(M,M+1)}) + 4\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(M+1,M+2)}). \end{aligned}$$

As $\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(-1,0)}) < \mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(-2,-1)})$ – because the kernel is radially decreasing and the interactions are negative by (6) – and $\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(M,M+1)}) \rightarrow 0$ as $M \uparrow \infty$ (along with the other similar terms), we deduce that there exists $M^*(s) > 0$ such that for any $M > M^*(s)$

$$(13) \quad \mathcal{E}_s(u, u) < \mathcal{E}_s(u^*, u^*).$$

Remark 4.2. A close inspection of $\mathcal{E}_s(u, u) - \mathcal{E}_s(u^*, u^*)$ can actually show that $M^*(s) = 1$ for any $s \in (1, 3/2)$. Indeed, let us define the functions

$$f: [1, \infty) \longrightarrow \mathbb{R} \qquad g: [1, \infty) \longrightarrow \mathbb{R}$$

$$r \longmapsto -2\mathcal{E}_{s-1}(\mathbf{1}_{(0,1)}, \mathbf{1}_{(r,r+1)}) \qquad \text{and} \qquad r \longmapsto f(r+1) - f(r).$$

As f is positive and decreasing, g is negative and

$$\mathcal{E}_s(u, u) - \mathcal{E}_s(u^*, u^*) = g(1) - 2g(M) + g(2M - 1).$$

Note that the following identities hold for $r > 1$:

$$f(r) = c_{1,s-1} \int_0^1 \int_0^1 (y - x + r)^{1-2s} dy dx,$$

$$f'''(r) = c_{1,s-1}(1 - 2s)(-2s)(-1 - 2s) \int_0^1 \int_0^1 (y - x + r)^{-2-2s} dy dx < 0$$

so that

$$g''(r) = f''(r+1) - f''(r) < 0 \quad \text{for } r > 1.$$

Using twice a second-order Taylor expansion centred at M with the remainder in Lagrange form, one deduces

$$\mathcal{E}_s(u, u) - \mathcal{E}_s(u^*, u^*) = g(1) - 2g(M) + g(2M - 1) = g''(\xi)(M - 1)^2 < 0.$$

Remark 4.3. In view of (12) and (13), it follows immediately that a Pólya–Szegő inequality or a reversed version of it cannot hold.

5. Analysis of the two-ball domain

In this section, let $D_1, D_2 \subset \mathbb{R}^n$ be two open balls with $D_1 \cap D_2 = \emptyset$ and $\Omega = D_1 \cup D_2$. We perform a spectral analysis of the operator $(-\Delta)^s$ on the space $\mathcal{H}_0^s(\Omega)$, by showing what we announced in the introduction, *i.e.*, that the first eigenvalue is positive and simple, with an associated eigenfunction which is possibly sign-changing according to the value of s : this is done in Subsection 5.1. Moreover, we improve the analysis carried out in [3] and [2] by giving sufficient conditions under which maximum principles are recovered on Ω : this is done in Subsection 5.2.

5.1. The first eigenfunction.

Lemma 5.1. *For $U \subset \mathbb{R}^n$ open and bounded, denote by*

$$\lambda(U) := \min_{\substack{u \in \mathcal{H}_0^s(U) \\ u \neq 0}} \frac{\mathcal{E}_s(u, u)}{\|u\|_{L^2(U)}^2}$$

the first Dirichlet eigenvalue of $(-\Delta)^s$ in U . Then $\lambda(U) > 0$ and it is attained by some $\varphi \in \mathcal{H}_0^s(U)$.

The proof of this lemma is standard and we omit it here.

In the following, we denote by π_i^\pm , $i = 1, 2$, the projections of $\mathcal{H}_0^s(D_i)$ onto their positive and negative cones, as defined and analysed in Appendix A.

Proposition 5.2. *Let $[s] \in 2\mathbb{N}_0 + 1$. The function $\tilde{\varphi} = 2\pi_1^+ \varphi_1 - \varphi_1 + 2\pi_2^- \varphi_2 - \varphi_2$ satisfies*

$$(14) \qquad \|\tilde{\varphi}\|_{L^2(\Omega)} \geq \|\varphi\|_{L^2(\Omega)}, \quad \mathcal{E}_s(\tilde{\varphi}, \tilde{\varphi}) \leq \mathcal{E}_s(\varphi, \varphi).$$

If instead $[s] \in 2\mathbb{N}_0$, then $\tilde{\varphi} = 2\pi_1^+ \varphi_1 - \varphi_1 + 2\pi_2^+ \varphi_2 - \varphi_2$ satisfies (14).

Proof: Let us suppose $[s] \in 2\mathbb{N}_0 + 1$; the case $[s] \in 2\mathbb{N}_0$ will follow under suitable minor modifications.

By the maximum principle on balls [1, Theorem 1.1] and (23), it holds that $\pi_1^+ \varphi_1 \geq \varphi_1$ and $\pi_2^- \varphi_2 \leq \varphi_2$. One has

$|\varphi| \leq |\varphi_1 - \pi_1^+ \varphi_1 + \pi_1^+ \varphi_1| + |\varphi_2 - \pi_2^- \varphi_2 + \pi_2^- \varphi_2| = 2\pi_1^+ \varphi_1 - \varphi_1 - 2\pi_2^- \varphi_2 + \varphi_2 = |\tilde{\varphi}|$
and this proves the claim about the L^2 norm. Also,

$$\begin{aligned}
 \mathcal{E}_s(\tilde{\varphi}, \tilde{\varphi}) &= \mathcal{E}_s(2\pi_1^+ \varphi_1 - \varphi_1 + 2\pi_2^- \varphi_2 - \varphi_2, 2\pi_1^+ \varphi_1 - \varphi_1 + 2\pi_2^- \varphi_2 - \varphi_2) \\
 &= \mathcal{E}_s(2\pi_1^+ \varphi_1 - \varphi_1, 2\pi_1^+ \varphi_1 - \varphi_1) + \mathcal{E}_s(2\pi_2^- \varphi_2 - \varphi_2, 2\pi_2^- \varphi_2 - \varphi_2) \\
 &\quad + 2\mathcal{E}_s(2\pi_1^+ \varphi_1 - \varphi_1, 2\pi_2^- \varphi_2 - \varphi_2) \\
 (15) \quad &\leq \mathcal{E}_s(\varphi_1, \varphi_1) + \mathcal{E}_s(\varphi_2, \varphi_2) - 2\mathcal{E}_s(\varphi_1, \varphi_2) \\
 &\quad + 4\mathcal{E}_s(\pi_1^+ \varphi_1 - \varphi_1, \pi_2^- \varphi_2) + 4\mathcal{E}_s(\pi_1^+ \varphi_1, \pi_2^- \varphi_2 - \varphi_2) \\
 &\leq \mathcal{E}_s(\varphi, \varphi),
 \end{aligned}$$

where we have used (28) and the inner product of functions with segregated supports (6). □

Lemma 5.3. $\lambda(\Omega) < \min\{\lambda(D_1), \lambda(D_2)\}$. As a consequence, $\varphi \not\equiv 0$ in D_1 and $\varphi \not\equiv 0$ in D_2 .

Proof: Suppose, without loss of generality, that $\lambda(D_1) \leq \lambda(D_2)$. Denote by $u_1 \in \mathcal{H}_0^s(D_1)$ and $u_2 \in \mathcal{H}^s(D_2)$ the normalized (positive) eigenfunctions associated to $(-\Delta)^s$ on D_1 and D_2 respectively. For any $\alpha \in [0, 1]$:

$$\begin{aligned}
 \mathcal{E}_s(\sqrt{1 - \alpha^2} u_1 + (-1)^{[s]} \alpha u_2, \sqrt{1 - \alpha^2} u_1 + (-1)^{[s]} \alpha u_2) \\
 = (1 - \alpha^2)\lambda(D_1) + \alpha^2\lambda(D_2) + 2(-1)^{[s]} \alpha \sqrt{1 - \alpha^2} \mathcal{E}_s(u_1, u_2).
 \end{aligned}$$

When $\alpha > 0$ is very small, one has $\alpha^2\lambda(D_2) + 2(-1)^{[s]} \alpha \sqrt{1 - \alpha^2} \mathcal{E}_s(u_1, u_2) < 0$ and therefore there is a suitable choice $\alpha_* \in (0, 1)$ such that

$$\mathcal{E}_s(\sqrt{1 - \alpha_*^2} u_1 - \alpha_* u_2, \sqrt{1 - \alpha_*^2} u_1 - \alpha_* u_2) < \lambda(D_1) \leq \lambda(D_2),$$

which proves the first claim.

If it were $\varphi \equiv 0$ in D_2 , then it would also hold that

$$\varphi \in \mathcal{H}_0^s(D_1), \quad \mathcal{E}_s(\varphi, \varphi) = \lambda(\Omega) < \lambda(D_1),$$

which contradicts the definition of $\lambda(D_1)$. □

Theorem 5.4. $\lambda(\Omega)$ is simple.

If $[s] \in 2\mathbb{N}_0 + 1$, then φ is of one sign on D_1 and of the opposite one in D_2 .

If instead $[s] \in 2\mathbb{N}_0$, then φ is of constant sign in Ω .

Proof: Let us suppose $[s] \in 2\mathbb{N}_0 + 1$; the case $[s] \in 2\mathbb{N}_0$ will follow after suitable minor modifications.

In the notations of Proposition 5.2, we already know that

$$\mathcal{E}_s(\tilde{\varphi}, \tilde{\varphi}) \leq \mathcal{E}_s(\varphi, \varphi)$$

and, in view of (15), equality can hold only if

$$\mathcal{E}_s(\pi_1^+ \varphi_1, \pi_2^- \varphi_2 - \varphi_2) + \mathcal{E}_s(\pi_1^+ \varphi_1 - \varphi_1, \pi_2^- \varphi_2) = 0.$$

This last condition in turn holds if and only if one of the following is true:

- (i) $\pi_1^+ \varphi_1 = 0$ and $\pi_2^- \varphi_2 = 0$, which means that $\varphi_1 \leq 0$ and $\varphi_2 \geq 0$;
- (ii) $\pi_2^- \varphi_2 = \varphi_2$ and $\pi_1^+ \varphi_1 = \varphi_1$, which means that $\varphi_2 \leq 0$ and $\varphi_1 \geq 0$;

(iii) $\varphi_1 = 0$ or $\varphi_2 = 0$, but in this case the minimizing condition would be broken; see Lemma 5.3.

This proves the claim about the sign of the eigenfunction.

Suppose now that we have two minimizers φ and ψ . Without loss of generality (by the previous claim, they are necessarily of one sign on each ball) we can suppose that $\varphi\psi \leq 0$ in Ω . Define, for any $t \in [0, 1]$, $w_t := t\varphi + \sqrt{1 - t^2}\psi$. One can verify that w_t is either still a minimizer or trivial. In the first case, w_t must be of one sign on each ball and, since $w_1 = \varphi$ and $w_0 = \psi$ are of opposite sign throughout Ω , there must exist a $t_1 \in (0, 1)$ such that $w_{t_1} = 0$ in D_1 . This feature is not compatible with minimality (compare with (iii) above) so we must conclude $w_{t_1} = 0$ in Ω . Thus φ and ψ are not linearly independent. \square

5.2. Partial recovery of maximum principles.

Proposition 5.5. *Suppose $|D_1| = |D_2| = \omega_n$ and $\text{dist}(D_1, D_2) \geq 2$. Denote by $\tau: \mathbb{R}^n \rightarrow \mathbb{R}^n$ the inversion of \mathbb{R}^n such that $\tau(D_1) = D_2$. Let $g \in L^2(D_1)$, $g \geq 0$, and $f \in L^2(\Omega)$, $f = g - (-1)^{\lfloor s \rfloor} g \circ \tau$; then the weak solution $u \in \mathcal{H}_0^s(\Omega)$ of*

$$(-\Delta)^s u = f \quad \text{in } \Omega$$

is of the form $u = v - (-1)^{\lfloor s \rfloor} v \circ \tau$, $v \in \mathcal{H}_0^s(D_1)$, with $v \geq 0$ in D_1 .

Proof: The solution $u \in \mathcal{H}_0^s(\Omega)$ is given by the Green representation

$$u(x) = \int_{\Omega} G_{\Omega}(x, y) f(y) dy, \quad x \in \Omega.$$

Using the symmetry of f , u can be alternatively written

$$u(x) = \int_{D_1} [G_{\Omega}(x, y) - (-1)^{\lfloor s \rfloor} G_{\Omega}(x, \tau(y))] g(y) dy, \quad x \in \Omega.$$

We claim that

$$(16) \quad G_{\Omega}(x, y) - (-1)^{\lfloor s \rfloor} G_{\Omega}(x, \tau(y)) \geq 0, \quad x, y \in D_1.$$

The Green function $G_{\Omega}(x, \cdot)$ can be represented, for $x \in D_1$, by making use of the Poisson kernel (see [2, Proposition 3.4]) as

$$G_{\Omega}(x, y) = \begin{cases} G_{D_1}(x, y) + \int_{D_2} \Gamma_{D_1}(y, z) G_{\Omega}(x, z) dz, & y \in D_1, \\ \int_{D_1} \Gamma_{D_2}(y, z) G_{\Omega}(x, z) dz = \int_{D_2} \Gamma_{D_1}(\tau(y), z) G_{\Omega}(x, \tau(z)) dz, & y \in D_2, \end{cases}$$

and thus

$$G_{\Omega}(x, y) - (-1)^{\lfloor s \rfloor} G_{\Omega}(x, \tau(y)) = G_{D_1}(x, y) + \int_{D_2} \Gamma_{D_1}(y, z) G_{\Omega}(x, z) dz - (-1)^{\lfloor s \rfloor} \int_{D_2} \Gamma_{D_1}(y, z) G_{\Omega}(x, \tau(z)) dz, \quad x, y \in D_1.$$

Denoting $w_x(y) := G_{\Omega}(x, y) - (-1)^{\lfloor s \rfloor} G_{\Omega}(x, \tau(y))$, $x \in D_1$, $y \in \Omega$, we have for $y \in D_1$

$$\begin{aligned} w_x(y) &= G_{D_1}(x, y) + \int_{D_2} \Gamma_{D_1}(y, z) w_x(z) dz \\ &= G_{D_1}(x, y) + \int_{D_2} \Gamma_{D_1}(y, z) G_{D_1}(x, \tau(z)) dz \\ &\quad + \int_{D_2} \left(\int_{D_2} \Gamma_{D_1}(y, z) \Gamma_{D_1}(z, \xi) dz \right) w_x(\xi) d\xi. \end{aligned}$$

The rest of the proof consists of two different steps: the first one is to prove

$$(17) \quad G_{D_1}(x, y) + \int_{D_2} \Gamma_{D_1}(y, z) G_{D_1}(x, \tau(z)) dz \geq 0, \quad x, y \in D_1,$$

when $\text{dist}(D_1, D_2) \geq 2$; once this is done, the second one is to bootstrap the resulting inequality

$$(18) \quad w_x(y) \geq \int_{D_2} \left(\int_{D_2} \Gamma_{D_1}(y, z) \Gamma_{D_1}(z, \xi) dz \right) w_x(\xi) d\xi$$

to get (16) in the end, as desired.

As to (17), we use the explicit expression of G_{D_1} known as *Boggio's formula* (see [1, Theorem 1.1], [12, Theorem 1], or also [15, Remark 1]),

$$G_{D_1}(x, y) = k_{n,s} \frac{(1 - |x|^2)^s (1 - |y|^2)^s}{|x - y|^n} \int_0^1 \frac{\eta^{s-1}}{(\rho(x, y)\eta + 1)^{n/2}} d\eta, \quad x, y \in D_1,$$

$$\rho(x, y) = \frac{(1 - |x|^2)(1 - |y|^2)}{|x - y|^2},$$

$$k_{n,s} = \frac{\Gamma(n/2)}{2^{2s} \pi^{n/2} \Gamma(s)^2},$$

and the explicit formula for the Poisson kernel Γ_{D_1} (see [2, Theorem 1.1])

$$\Gamma_{D_1}(y, z) = \frac{(-1)^{\lfloor s \rfloor} \gamma_{n,s}}{|y - z|^n} \left(\frac{1 - |y|^2}{|z|^2 - 1} \right)^s, \quad y \in D_1, z \in \mathbb{R}^n \setminus \bar{D}_1,$$

$$\gamma_{n,s} = \frac{\Gamma(n/2)}{\pi^{n/2} \Gamma(s - \lfloor s \rfloor) \Gamma(1 - s + \lfloor s \rfloor)}.$$

In order to get (17) we estimate

$$k_{n,s} \frac{(1 - |x|^2)^s}{|x - y|^n} \int_0^1 \frac{\eta^{s-1}}{(\rho(x, y)\eta + 1)^{n/2}} d\eta - \gamma_{n,s} \int_{D_2} \frac{G_{D_1}(x, \tau(z))}{|y - z|^n (|z|^2 - 1)^s} dz$$

$$\geq \frac{k_{n,s}}{s} \frac{(1 - |x|^2)^s}{((1 - |x|^2)(1 - |y|^2) + |x - y|^2)^{n/2}} - \frac{\gamma_{n,s}}{2^n} \int_{D_2} G_{D_1}(x, \tau(z)) dz,$$

where we have used $\text{dist}(D_1, D_2) \geq 2$. Next,

$$\int_{D_2} G_{D_1}(x, \tau(z)) dz = \int_{D_1} G_{D_1}(x, z) dz$$

gives rise to the torsion function of D_1 (see [30, Lemma 2.1])

$$\int_{D_1} G_{D_1}(x, z) dz = \frac{\Gamma(n/2)}{2^{2s} \Gamma(n/2 + s) \Gamma(1 + s)} (1 - |x|^2)^s$$

so that

$$k_{n,s} \frac{(1 - |x|^2)^s}{|x - y|^n} \int_0^1 \frac{\eta^{s-1}}{(\rho(x, y)\eta + 1)^{n/2}} d\eta - \gamma_{n,s} \int_{D_2} \frac{G_{D_1}(x, \tau(z))}{|y - z|^n (|z|^2 - 1)^s} dz$$

$$\geq \frac{k_{n,s}}{s 2^n} (1 - |x|^2)^s - \frac{\gamma_{n,s}}{2^{2s}} \frac{\Gamma(n/2)}{2^{2s} \Gamma(n/2 + s) \Gamma(1 + s)} (1 - |x|^2)^s,$$

where we have used $\text{dist}(D_1, D_2) \geq 2$ and $(1 - |x|^2)(1 - |y|^2) + |x - y|^2 \leq 4$ for $x, y \in D_1$. The above estimate provides (17) if we check

$$\frac{1}{s\Gamma(s)^2} - \frac{1}{\Gamma(s - \lfloor s \rfloor)\Gamma(1 - s + \lfloor s \rfloor)} \frac{\Gamma(n/2)}{\Gamma(n/2 + s)\Gamma(1 + s)} \geq 0,$$

which, by using standard properties² of the Γ function, is equivalent to

$$(19) \quad \frac{\sin((s - \lfloor s \rfloor)\pi)}{\pi} \frac{\Gamma(n/2)\Gamma(s)}{\Gamma(n/2 + s)} \leq 1.$$

We now use the integral representation of the *Beta function* in order to deduce

$$\frac{\Gamma(n/2)\Gamma(s)}{\Gamma(n/2 + s)} = \int_0^1 r^{n/2-1}(1-r)^{s-1} dr \leq \int_0^1 r^{n/2-1} dr = \frac{2}{n} \leq 2.$$

This proves (19) and, therefore, (17).

From the above, (18) follows and we can proceed as in [2] to conclude the proof. \square

Theorem 5.6. *Under the assumptions of Proposition 5.5, the weak solution $v \in \mathcal{H}_0^s(D_1)$ of*

$$(-\Delta)^s v - (-1)^{\lfloor s \rfloor} (-\Delta)^s (v \circ \tau) = g \quad \text{in } D_1$$

is nonnegative.

Proof: It follows immediately from Proposition 5.5. \square

6. The Faber–Krahn inequality in dimension one

In this section we restrict to $n = 1$ and consider an open set $\Omega \subset \mathbb{R}$ with $|\Omega| = 2$. We denote

$$\Omega = \bigcup_{k \in \mathcal{M}} I_k, \quad \text{where } I_k \text{'s are disjoint open intervals, } \mathcal{M} \subseteq \mathbb{N},$$

$$\varphi = \sum_{k \in \mathcal{M}} \varphi_k, \quad \varphi_k \in \mathcal{H}_0^s(I_k), \text{ for any } k \in \mathcal{M}.$$

Lemma 6.1. *For any $s > 0$ and $k \in \mathcal{M}$, any first eigenfunction φ is of one sign in I_k .*

Proof: Denote by π_k^+ the projection onto the positive cone $\mathcal{C}^+(I_k)$ as defined in (20)-(21). Consider the function

$$\tilde{\varphi} = \varphi - 2\varphi_k + 2\pi_k^+ \varphi_k \in \mathcal{H}_0^s(\Omega)$$

and compute

$$\|\tilde{\varphi}\|_{L^2(\Omega)} = \|\varphi - \varphi_k\|_{L^2(\Omega \setminus I_k)} + \|2\pi_k^+ \varphi_k - \varphi_k\|_{L^2(I_k)}.$$

If φ_k changes sign, then $|\varphi_k| = |\varphi_k - \pi_k^+ \varphi_k + \pi_k^+ \varphi_k| < 2\pi_k^+ \varphi_k - \varphi_k$ in I_k by the strong maximum principles³ on intervals and (23). Therefore $\|\varphi\|_{L^2(\Omega)} < \|\tilde{\varphi}\|_{L^2(\Omega)}$.

Furthermore,

$$\begin{aligned} \mathcal{E}_s(\tilde{\varphi}, \tilde{\varphi}) &= \mathcal{E}_s(\varphi - \varphi_k, \varphi - \varphi_k) + \mathcal{E}_s(2\pi_k^+ \varphi_k - \varphi_k, 2\pi_k^+ \varphi_k - \varphi_k) + 2\mathcal{E}_s(2\pi_k^+ \varphi_k - \varphi_k, \varphi - \varphi_k) \\ &\leq \mathcal{E}_s(\varphi - \varphi_k, \varphi - \varphi_k) + \mathcal{E}_s(\varphi_k, \varphi_k) + 2\mathcal{E}_s(2\pi_k^+ \varphi_k - \varphi_k, \varphi - \varphi_k) \\ &= \mathcal{E}_s(\varphi, \varphi) + 4\mathcal{E}_s(\pi_k^+ \varphi_k - \varphi_k, \varphi - \varphi_k) \end{aligned}$$

by (28). Up to changing the sign of $\varphi - \varphi_k$, we deduce $\mathcal{E}_s(\tilde{\varphi}, \tilde{\varphi}) \leq \mathcal{E}_s(\varphi, \varphi)$ by (6). \square

²We need in particular the *Euler's reflection formula*: $\Gamma(\alpha)\Gamma(1 - \alpha) = \frac{\pi}{\sin(\pi\alpha)}$ for $\alpha \in \mathbb{C} \setminus \mathbb{Z}$.

³This follows by the positivity of the Green function on balls; see [1, 12].

Lemma 6.2. *For any $s > 0$, if there exist $k, h \in \mathcal{M}, k \neq h$, such that $\text{dist}(I_k, I_h) = 0$, then there exists $x_0 \in \mathbb{R}$ such that $I_k \cup \{x_0\} \cup I_h$ is connected and $\lambda(\Omega \cup \{x_0\}) \leq \lambda(\Omega)$.*

Proof: This simply follows by the inclusion $\mathcal{H}_0^s(\Omega) \subseteq \mathcal{H}_0^s(\Omega \cup \{x_0\})$. □

6.1. When the integer part of s is even.

Proposition 6.3. *Let $\lfloor s \rfloor \in 2\mathbb{N}_0$. Then $\varphi \geq 0$ in Ω .*

Proof: By Lemma 6.1 we already know that φ is of one sign on each connected component of Ω , so that $\varphi^+, \varphi^- \in \mathcal{H}_0^s(\Omega)$. As $\lfloor s \rfloor \in 2\mathbb{N}_0$ by assumption, one has by (6)

$$\begin{aligned} \mathcal{E}_s(|\varphi|, |\varphi|) &= \mathcal{E}_s(\varphi^+, \varphi^+) + \mathcal{E}_s(\varphi^-, \varphi^-) + 2\mathcal{E}_s(\varphi^+, \varphi^-) \\ &< \mathcal{E}_s(\varphi^+, \varphi^+) + \mathcal{E}_s(\varphi^-, \varphi^-) - 2\mathcal{E}_s(\varphi^+, \varphi^-) = \mathcal{E}_s(\varphi, \varphi). \end{aligned} \quad \square$$

Theorem 6.4. *Let $\lfloor s \rfloor \in 2\mathbb{N}_0$. Then $\lambda(-1, 1) \leq \lambda(\Omega)$ for any open set $\Omega \subset \mathbb{R}$ with $|\Omega| = 2$. Moreover, equality holds if and only if Ω is an interval.*

Proof: By Lemma 6.2, we can restrict our attention to those Ω 's satisfying $\text{dist}(I_k, I_h) > 0$ for all $k, h \in \mathcal{M}, k \neq h$.

Let $x_0 \in \mathbb{R} \setminus \Omega$ be such that $(-\infty, x_0) \cap \Omega \neq \emptyset \neq \Omega \cap (x_0, +\infty)$. Write $\varphi = \varphi \mathbf{1}_{(-\infty, x_0)} + \varphi \mathbf{1}_{(x_0, +\infty)}$. We have

$$\begin{aligned} \mathcal{E}_s(\varphi, \varphi) &= \mathcal{E}_s(\varphi \mathbf{1}_{(-\infty, x_0)}, \varphi \mathbf{1}_{(-\infty, x_0)}) + \mathcal{E}_s(\varphi \mathbf{1}_{(x_0, +\infty)}, \varphi \mathbf{1}_{(x_0, +\infty)}) \\ &\quad + 2\mathcal{E}_s(\varphi \mathbf{1}_{(-\infty, x_0)}, \varphi \mathbf{1}_{(x_0, +\infty)}). \end{aligned}$$

By Proposition 6.3 and the fact that $\lfloor s \rfloor \in 2\mathbb{N}_0$, the last term is nonpositive; indeed

$$\mathcal{E}_s(\varphi \mathbf{1}_{(-\infty, x_0)}, \varphi \mathbf{1}_{(x_0, +\infty)}) = -\frac{c_{1,s}}{2} \int_{-\infty}^{x_0} \int_{x_0}^{+\infty} \frac{\varphi(x) \varphi(y)}{(y-x)^{1+2s}} dy dx.$$

Since there exists $\varepsilon > 0$ such that $(x_0 - \varepsilon, x_0) \cap \Omega = \emptyset$, then substituting $\varphi(y) \mathbf{1}_{(x_0, +\infty)}(y)$ with $\varphi(y + \varepsilon) \mathbf{1}_{(x_0, +\infty)}(y)$ gives

$$\begin{aligned} - \int_{-\infty}^{x_0} \int_{x_0 - \varepsilon}^{+\infty} \frac{\varphi(x) \varphi(y + \varepsilon)}{(y-x)^{1+2s}} dy dx &= - \int_{-\infty}^{x_0} \int_{x_0}^{+\infty} \frac{\varphi(x) \varphi(y)}{(y-x-\varepsilon)^{1+2s}} dy dx \\ &\leq - \int_{-\infty}^{x_0} \int_{x_0}^{+\infty} \frac{\varphi(x) \varphi(y)}{(y-x)^{1+2s}} dy dx. \end{aligned}$$

We then deduce that $\lambda(\Omega)$ is minimized when Ω is connected.

We now prove the second claim in the statement. In order to do so, we only need to prove that, for any $x_0 \in (-1, 1)$, the following holds:

$$\lambda(-1, 1) < \lambda((-1, 1) \setminus \{x_0\}).$$

This is a direct consequence of the strong maximum principle in $(-1, 1)$, which can be stated as a corollary of the positivity of the Green function; see [1, Theorem 1.1]. □

6.2. When the integer part of s is odd.

Proposition 6.5. *Let $\lfloor s \rfloor \in 2\mathbb{N}_0 + 1$. Suppose that $\Omega \subset \mathbb{R}$ is an open set with $|\Omega| = 2$ and that \mathcal{M} is finite. Then $\lambda(-1, 1) \leq \lambda(\Omega)$ and equality holds if and only if Ω is an interval.*

Proof: Suppose that Ω is not connected: by Lemma 6.2, we can also suppose that

$$\delta_0 = \min\{\text{dist}(I_k, I_h) : k, h \in \mathcal{M}, k \neq h\} > 0.$$

In this case, φ changes sign in Ω : this simply follows from Lemma 6.1 and (6) and, moreover, $\varphi^+, \varphi^- \in \mathcal{H}_0^s(\Omega)$. Consider the function

$$u_\delta(x) = \varphi^+(x) - \varphi^-(x - \delta), \quad x \in \mathbb{R},$$

where δ ranges in the largest open interval I such that $\{x \in \mathbb{R} : \varphi^+(x) > 0\} \cap \{x \in \mathbb{R} : \varphi^-(x - \delta) > 0\} = \emptyset$ for any $\delta \in I$: this interval is not empty because, by assumption, contains at least $(-\delta_0, \delta_0)$. It holds that $u_\delta \in H^s(\mathbb{R})$ and, moreover, $\|u_\delta\|_{L^2(\mathbb{R})} = \|\varphi\|_{L^2(\mathbb{R})}$. Also,

$$\begin{aligned} \mathcal{E}_s(u_\delta, u_\delta) &= \mathcal{E}_s(\varphi^+, \varphi^+) + \mathcal{E}_s(\varphi^-, \varphi^-) - c_{1,s} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{\varphi^+(x) \varphi^-(y - \delta)}{|x - y|^{1+2s}} dx dy \\ &= \mathcal{E}_s(\varphi^+, \varphi^+) + \mathcal{E}_s(\varphi^-, \varphi^-) - c_{1,s} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{\varphi^+(x) \varphi^-(y)}{|x - y - \delta|^{1+2s}} dx dy. \end{aligned}$$

We now differentiate twice in δ and obtain

$$\begin{aligned} \frac{1}{c_{1,s}} \frac{d^2}{d\delta^2} \mathcal{E}_s(u_\delta, u_\delta) &= -(1 + 2s) \frac{d}{d\delta} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{\varphi^+(x) \varphi^-(y)}{|x - y - \delta|^{2+2s}} \text{sgn}(x - y - \delta) dx dy \\ &= -(1 + 2s) \frac{d}{d\delta} \int_{\mathbb{R}} \int_y^{+\infty} \frac{\varphi^+(x) \varphi^-(y)}{(x - y - \delta)^{2+2s}} dx dy \\ &\quad + (1 + 2s) \frac{d}{d\delta} \int_{\mathbb{R}} \int_{-\infty}^y \frac{\varphi^+(x) \varphi^-(y)}{(y + \delta - x)^{2+2s}} dx dy \\ &= -(1 + 2s)(2 + 2s) \int_{\mathbb{R}} \int_y^{+\infty} \frac{\varphi^+(x) \varphi^-(y)}{(x - y - \delta)^{3+2s}} dx dy \\ &\quad - (1 + 2s)(2 + 2s) \int_{\mathbb{R}} \int_{-\infty}^y \frac{\varphi^+(x) \varphi^-(y)}{|x - y - \delta|^{3+2s}} dx dy \\ &= -(1 + 2s)(2 + 2s) \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{\varphi^+(x) \varphi^-(y)}{|x - y - \delta|^{3+2s}} dx dy < 0, \quad \text{for any } \delta \in I. \end{aligned}$$

From this we deduce that $I \ni \delta \mapsto \mathcal{E}_s(u_\delta, u_\delta)$ is minimized at one of the two extremal points of I : this configuration cannot be minimal again for Lemma 6.2. So, the first eigenvalue of Ω would be larger than the one of a domain with one less connected component. The conclusion follows by iterating the above (recall that \mathcal{M} is finite by assumption). \square

Theorem 6.6. *Let $[s] \in 2\mathbb{N}_0 + 1$. Then $\lambda(-1, 1) \leq \lambda(\Omega)$ for any open set $\Omega \subset \mathbb{R}$ with $|\Omega| = 2$.*

Proof: By Proposition 6.5, we only need to prove that the claim holds when \mathcal{M} is countable. We proceed via a continuity argument. Write

$$\begin{aligned} \Omega &= \bigcup_{k \in \mathbb{N}} I_k, \quad \text{where } I_k \text{'s are disjoint open intervals,} & \Omega_N &= \bigcup_{k=1}^N I_k, \quad \text{for any } N \in \mathbb{N}, \\ \varphi &= \sum_{k \in \mathbb{N}} \varphi_k, \quad \varphi_k \in \mathcal{H}_0^s(I_k), \quad \text{for any } k \in \mathbb{N}, & u_N &= \sum_{k=1}^N \varphi_k, \quad \text{for any } N \in \mathbb{N}. \end{aligned}$$

By the inclusion $\Omega_N \subset \Omega$, one simply has $\lambda(\Omega_N) \geq \lambda(\Omega)$. Moreover,

$$\begin{aligned} \lambda(\Omega) &= \mathcal{E}_s(\varphi, \varphi) = \mathcal{E}_s(u_N, u_N) + 2\mathcal{E}_s(\varphi - u_N, u_N) + \mathcal{E}_s(\varphi - u_N, \varphi - u_N) \\ &\geq \lambda(\Omega_N) \|u_N\|_{L^2(\mathbb{R}^n)}^2 + 2\mathcal{E}_s(\varphi - u_N, u_N) + \mathcal{E}_s(\varphi - u_N, \varphi - u_N). \end{aligned}$$

By Proposition 6.5, we have that $\lambda(\Omega_N)$ is greater than the first eigenvalue of the ball with volume $|\Omega_N|$, and then by the scaling properties of $(-\Delta)^s$

$$\lambda(\Omega_N) \geq \left(\frac{2}{|\Omega_N|} \right)^{2s} \lambda(-1, 1).$$

We have $u_N \rightarrow \varphi$ in $H^s(\mathbb{R}^n)$: indeed, by (5) and since $\varphi - u_N$ is supported in $\Omega \setminus \Omega_N$,

$$\begin{aligned} \mathcal{E}_s(\varphi - u_N, \varphi - u_N) &= \mathcal{E}_{s-[s]}(D(\varphi - u_N), D(\varphi - u_N)), \quad \text{where } D = \nabla(-\Delta)^{([s]-1)/2} \\ &= \frac{c_{1,s-[s]}}{2} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{|D(\varphi - u_N)(x) - D(\varphi - u_N)(y)|^2}{|x - y|^{1+2s-2[s]}} dx dy \\ &= \frac{c_{1,s-[s]}}{2} \int_{\Omega \setminus \Omega_N} \int_{\Omega \setminus \Omega_N} \frac{|D\varphi(x) - D\varphi(y)|^2}{|x - y|^{1+2s-2[s]}} dx dy \\ &\quad + c_{1,s-[s]} \int_{\mathbb{R} \setminus (\Omega \setminus \Omega_N)} \int_{\Omega \setminus \Omega_N} \frac{|D\varphi(x)|^2}{|x - y|^{1+2s-2[s]}} dx dy \\ &\leq \frac{c_{1,s-[s]}}{2} \int_{\Omega \setminus \Omega_N} \int_{\Omega \setminus \Omega_N} \frac{|D\varphi(x) - D\varphi(y)|^2}{|x - y|^{1+2s-2[s]}} dx dy \\ &\quad + c_{1,s-[s]} \int_{\Omega \setminus \Omega_N} |D\varphi(x)|^2 \text{dist}(x, \mathbb{R} \setminus (\Omega \setminus \Omega_N))^{-2s+2[s]} dx, \end{aligned}$$

which converges to 0 as $N \uparrow \infty$ and $\Omega_N \nearrow \Omega$ by, for example, [38, Section 4.3.2, equation (7)]. We deduce

$$\begin{aligned} \lambda(\Omega) &\geq \lim_{N \uparrow \infty} [\lambda(\Omega_N) \|u_N\|_{L^2(\mathbb{R}^n)}^2 + 2\mathcal{E}_s(\varphi - u_N, u_N) + \mathcal{E}_s(\varphi - u_N, \varphi - u_N)] \\ &\geq \lim_{N \uparrow \infty} \left(\frac{2}{|\Omega_N|} \right)^{2s} \|u_N\|_{L^2(\mathbb{R}^n)}^2 \lambda(-1, 1) = \lambda(-1, 1). \quad \square \end{aligned}$$

Appendix A. Projections onto the positive cone

Denote

$$(20) \quad \mathcal{C}^+(\Omega) := \{v \in \mathcal{H}_0^s(\Omega) : v \geq 0\}$$

and note how any $\varphi \in C_c^\infty(\Omega)$, $\varphi \geq 0$, lies inside $\mathcal{C}^+(\Omega)$. Define the map

$$\pi^+ : \mathcal{H}_0^s(\Omega) \longrightarrow \mathcal{C}^+(\Omega)$$

which associates to w its projection on $\mathcal{C}^+(\Omega)$, that is

$$(21) \quad \mathcal{E}_s(w - \pi^+ w, w - \pi^+ w) \leq \mathcal{E}_s(w - v, w - v), \quad \text{for any } v \in \mathcal{C}^+(\Omega).$$

We list below a number of inequalities that are needed throughout the paper:

$$(22) \quad \mathcal{E}_s(w - \pi^+w, v - \pi^+w) \leq 0, \quad \text{for any } v \in \mathcal{C}^+(\Omega),$$

$$(23) \quad 0 \leq \mathcal{E}_s(\pi^+w - w, v), \quad \text{for any } v \in \mathcal{C}^+(\Omega),$$

$$(24) \quad \mathcal{E}_s(\pi^+w - w, \pi^+w) = 0,$$

$$(25) \quad \mathcal{E}_s(w, \pi^+w) \leq \mathcal{E}_s(w, w),$$

$$(26) \quad \mathcal{E}_s(w, v) \leq \mathcal{E}_s(\pi^+w, v), \quad \text{for any } v \in \mathcal{C}^+(\Omega),$$

$$(27) \quad \mathcal{E}_s(\pi^+w, \pi^+w) \leq \mathcal{E}_s(w, w),$$

$$(28) \quad \mathcal{E}_s(2\pi^+w - w, 2\pi^+w - w) \leq \mathcal{E}_s(w, w).$$

Proof of (22): By convexity of $\mathcal{C}^+(\Omega)$

$$\pi^+w + t(v - \pi^+w) \in \mathcal{C}^+(\Omega), \quad \text{for any } t \in [0, 1], v \in \mathcal{C}^+(\Omega).$$

So,

$$\begin{aligned} \mathcal{E}_s(w - \pi^+w, w - \pi^+w) &\leq \mathcal{E}_s(w - \pi^+w - t(v - \pi^+w), w - \pi^+w - t(v - \pi^+w)) \\ &= \mathcal{E}_s(w - \pi^+w, w - \pi^+w) - 2t\mathcal{E}_s(w - \pi^+w, v - \pi^+w) + t^2\mathcal{E}_s(v - \pi^+w, v - \pi^+w), \end{aligned}$$

which implies

$$2\mathcal{E}_s(w - \pi^+w, v - \pi^+w) \leq t\mathcal{E}_s(v - \pi^+w, v - \pi^+w), \quad \text{for any } t \in (0, 1]$$

and we deduce (22) by considering t arbitrarily small. □

Proof of (23): From (22) it follows that

$$\mathcal{E}_s(\pi^+w - w, v) \geq \mathcal{E}_s(\pi^+w - w, \pi^+w),$$

where the right-hand side is nonnegative again thanks to (22) by considering the particular case $v = 0$. □

Proof of (24): As mentioned above, inequality \leq follows from (22) by considering the particular case $v = 0$. Inequality \geq follows from (23) by considering the particular case $v = \pi^+w$. □

Proof of (25): Using again (24) we deduce

$$\mathcal{E}_s(w, \pi^+w - w) = \mathcal{E}_s(w - \pi^+w, \pi^+w - w) + \mathcal{E}_s(\pi^+w, \pi^+w - w) \leq 0. \quad \square$$

Proof of (26): Since both v and π^+w belong to $\mathcal{C}^+(\Omega)$, then so does $2\pi^+w + v$. Then, by (22),

$$\begin{aligned} 0 &\geq \mathcal{E}_s(w - \pi^+w, (2\pi^+w + v) - \pi^+w) = \mathcal{E}_s(w - \pi^+w, \pi^+w + v) \\ &= \mathcal{E}_s(w - \pi^+w, \pi^+w) + \mathcal{E}_s(w - \pi^+w, v) \end{aligned}$$

and since, by (24), it holds that $\mathcal{E}_s(w - \pi^+w, \pi^+w) \geq 0$, we deduce $\mathcal{E}_s(w - \pi^+w, v) \leq 0$. □

Proof of (27): Write (exploiting (24) in all inequalities below)

$$\begin{aligned} \mathcal{E}_s(\pi^+w, \pi^+w) - \mathcal{E}_s(w, w) &= \mathcal{E}_s(\pi^+w - w, \pi^+w) + \mathcal{E}_s(w, \pi^+w) - \mathcal{E}_s(w, w) \\ &\leq \mathcal{E}_s(w, \pi^+w - w) = \mathcal{E}_s(w - \pi^+w, \pi^+w - w) + \mathcal{E}_s(\pi^+w, \pi^+w - w) \leq 0. \quad \square \end{aligned}$$

Proof of (28): Write

$$\mathcal{E}_s(2\pi^+w - w, 2\pi^+w - w) = \mathcal{E}_s(w, w) + 4\mathcal{E}_s(\pi^+w - w, \pi^+w) \leq \mathcal{E}_s(w, w),$$

where we have used (24). □

References

- [1] N. ABATANGELO, S. JAROHS, AND A. SALDAÑA, Green function and Martin kernel for higher-order fractional Laplacians in balls, *Nonlinear Anal.* **175** (2018), 173–190. DOI: 10.1016/j.na.2018.05.019.
- [2] N. ABATANGELO, S. JAROHS, AND A. SALDAÑA, Integral representation of solutions to higher-order fractional Dirichlet problems on balls, *Commun. Contemp. Math.* **20(8)** (2018), 1850002, 36 pp. DOI: 10.1142/S0219199718500025.
- [3] N. ABATANGELO, S. JAROHS, AND A. SALDAÑA, On the loss of maximum principles for higher-order fractional Laplacians, *Proc. Amer. Math. Soc.* **146(11)** (2018), 4823–4835. DOI: 10.1090/proc/14165.
- [4] N. ABATANGELO, S. JAROHS, AND A. SALDAÑA, Positive powers of the Laplacian: from hypersingular integrals to boundary value problems, *Commun. Pure Appl. Anal.* **17(3)** (2018), 899–922. DOI: 10.3934/cpaa.2018045.
- [5] N. ABATANGELO, S. JAROHS, AND A. SALDAÑA, Fractional Laplacians on ellipsoids, *Math. Eng.* **3(5)** (2021), Paper no. 038, 34 pp. DOI: 10.3934/mine.2021038.
- [6] N. ABATANGELO AND E. VALDINOCI, Getting acquainted with the fractional Laplacian, in: *Contemporary Research in Elliptic PDEs and Related Topics*, Springer INdAM Ser. **33**, Springer, Cham, 2019, pp. 1–105. DOI: 10.1007/978-3-030-18921-1_1.
- [7] M. S. ASHBAUGH AND R. D. BENGURIA, On Rayleigh’s conjecture for the clamped plate and its generalization to three dimensions, *Duke Math. J.* **78(1)** (1995), 1–17. DOI: 10.1215/S0012-7094-95-07801-6.
- [8] A. BAERNSTEIN, II, A unified approach to symmetrization, in: *Partial Differential Equations of Elliptic Type* (Cortona, 1992), Sympos. Math. **XXXV**, Cambridge University Press, Cambridge, 1994, pp. 47–91.
- [9] G. BOURDAUD AND Y. MEYER, Fonctions qui opèrent sur les espaces de Sobolev, *J. Funct. Anal.* **97(2)** (1991), 351–360. DOI: 10.1016/0022-1236(91)90006-Q.
- [10] C. BUCUR AND E. VALDINOCI, *Nonlocal Diffusion and Applications*, Lect. Notes Unione Mat. Ital. **20**, Springer, [Cham]Unione Matematica Italiana, Bologna, 2016. DOI: 10.1007/978-3-319-28739-3.
- [11] E. DI NEZZA, G. PALATUCCI, AND E. VALDINOCI, Hitchhiker’s guide to the fractional Sobolev spaces, *Bull. Sci. Math.* **136(5)** (2012), 521–573. DOI: 10.1016/j.bulsci.2011.12.004.
- [12] S. DIPIERRO AND H.-C. GRUNAU, Boggio’s formula for fractional polyharmonic Dirichlet problems, *Ann. Mat. Pura Appl. (4)* **196(4)** (2017), 1327–1344. DOI: 10.1007/s10231-016-0618-z.
- [13] R. J. DUFFIN, On a question of Hadamard concerning super-biharmonic functions, *J. Math. Physics* **27** (1949), 253–258. DOI: 10.1002/sapm1948271253.
- [14] R. J. DUFFIN, The maximum principle and biharmonic functions, *J. Math. Anal. Appl.* **3(3)** (1961), 399–405. DOI: 10.1016/0022-247X(61)90066-X.
- [15] B. DYDA, A. KUZNETSOV, AND M. KWAŚNICKI, Eigenvalues of the fractional Laplace operator in the unit ball, *J. Lond. Math. Soc. (2)* **95(2)** (2017), 500–518. DOI: 10.1112/jlms.12024.
- [16] P. R. GARABEDIAN, A partial differential equation arising in conformal mapping, *Pacific J. Math.* **1(4)** (1951), 485–524. DOI: 10.2140/PJM.1951.1.485.
- [17] N. GAROFALO, Fractional thoughts, in: *New Developments in the Analysis of Nonlocal operators*, Contemp. Math. **723**, American Mathematical Society, Providence, RI, 2019, pp. 1–135. DOI: 10.1090/conm/723/14569.
- [18] F. GAZZOLA, H.-C. GRUNAU, AND G. SWEERS, *Polyharmonic Boundary Value Problems. Positivity Preserving and Nonlinear Higher Order Elliptic Equations in Bounded Domains*, Lecture Notes in Math. **1991**, Springer-Verlag, Berlin, 2010. DOI: 10.1007/978-3-642-12245-3.
- [19] A. GRECO AND S. JAROHS, Foliated Schwarz symmetry of solutions to a cooperative system of equations involving nonlocal operators, *J. Elliptic Parabol. Equ.* **8(1)** (2022), 383–417. DOI: 10.1007/s41808-022-00155-y.
- [20] P. GRISVARD, *Elliptic Problems in Nonsmooth Domains*, Reprint of the 1985 original, With a foreword by Susanne C. Brenner, Classics Appl. Math. **69**, Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 2011. DOI: 10.1137/1.9781611972030.
- [21] J. HADAMARD, Mémoire sur le problème d’analyse relatif à l’équilibre des plaques élastiques encastrées, *Mém. Sav. étrang.* **33** (1908), 1–128. Reprinted in: Mémoire sur le problème d’analyse relatif à l’équilibre des plaques élastiques encastrées, in: *Œuvres de Jacques Hadamard. Tomes I, II, III, IV*, Éditions du Centre National de la Recherche Scientifique, Paris, 1968, pp. 515–641.
- [22] S. JAROHS, Symmetry of solutions to nonlocal nonlinear boundary value problems in radial sets, *NoDEA Nonlinear Differential Equations Appl.* **23(3)** (2016), Art. 32, 22 pp. DOI: 10.1007/s00030-016-0386-x.

- [23] S. JAROHS AND T. WETH, Symmetry via antisymmetric maximum principles in nonlocal problems of variable order, *Ann. Mat. Pura Appl. (4)* **195(1)** (2016), 273–291. DOI: 10.1007/s10231-014-0462-y.
- [24] V. A. KOZLOV, V. A. KONDRAT'EV, AND V. G. MAZ'YA, On sign variability and the absence of “strong” zeros of solutions of elliptic equations (Russian), *Izv. Akad. Nauk SSSR Ser. Mat.* **53(2)** (1989), 328–344; translation in: *Math. USSR-Izv.* **34(2)** (1990), 337–353. DOI: 10.1070/IM1990v034n02ABEH000649.
- [25] O. LOPES AND M. MARIŞ, Symmetry of minimizers for some nonlocal variational problems, *J. Funct. Anal.* **254(2)** (2008), 535–592. DOI: 10.1016/j.jfa.2007.10.004.
- [26] R. MUSINA AND A. I. NAZAROV, A note on truncations in fractional Sobolev spaces, *Bull. Math. Sci.* **9(1)** (2019), 1950001, 7 pp. DOI: 10.1142/S1664360719500012.
- [27] N. S. NADIRASHVILI, Rayleigh’s conjecture on the principal frequency of the clamped plate, *Arch. Rational Mech. Anal.* **129(1)** (1995), 1–10. DOI: 10.1007/BF00375124.
- [28] M. NAKAI AND L. SARIO, Green’s function of the clamped punctured disk, *J. Austral. Math. Soc. Ser. B* **20(2)** (1977), 175–181. DOI: 10.1017/S0334270000001557.
- [29] Y. J. PARK, Logarithmic Sobolev trace inequality, *Proc. Amer. Math. Soc.* **132(7)** (2004), 2075–2083. DOI: 10.1090/S0002-9939-03-07329-5.
- [30] X. ROS-OTON AND J. SERRA, Local integration by parts and Pohozaev identities for higher order fractional Laplacians, *Discrete Contin. Dyn. Syst.* **35(5)** (2015), 2131–2150. DOI: 10.3934/dcds.2015.35.2131.
- [31] S. G. SAMKO, A. A. KILBAS, AND O. I. MARICHEV, *Fractional Integrals and Derivatives: Theory and Applications*, Edited and with a foreword by S. M. Nikol’skiĭ, Translated from the 1987 Russian original, Revised by the authors, Gordon and Breach Science Publishers, Yverdon, 1993.
- [32] J. B. SEIF, On the Green’s function for the biharmonic equation in an infinite wedge, *Trans. Amer. Math. Soc.* **182** (1973), 241–260. DOI: 10.2307/1996533.
- [33] H. S. SHAPIRO AND M. TEGMARK, An elementary proof that the biharmonic Green function of an eccentric ellipse changes sign, *SIAM Rev.* **36(1)** (1994), 99–101. DOI: 10.1137/1036005.
- [34] G. SWEERS, An elementary proof that the triharmonic Green function of an eccentric ellipse changes sign, *Arch. Math. (Basel)* **107(1)** (2016), 59–62. DOI: 10.1007/s00013-016-0909-z.
- [35] G. SWEERS, Correction to: An elementary proof that the triharmonic Green function of an eccentric ellipse changes sign, *Arch. Math. (Basel)* **112(2)** (2019), 223–224. DOI: 10.1007/s00013-018-1274-x.
- [36] G. TALENTI, Elliptic equations and rearrangements, *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4)* **3(4)** (1976), 697–718.
- [37] G. TALENTI, On the first eigenvalue of the clamped plate, *Ann. Mat. Pura Appl. (4)* **129** (1981), 265–280. DOI: 10.1007/BF01762146.
- [38] H. TRIEBEL, *Interpolation Theory, Function Spaces, Differential operators*, North-Holland Math. Library **18**, North-Holland Publishing Co., Amsterdam-New York, 1978.
- [39] J. VAN SCHAFTINGEN, Universal approximation of symmetrizations by polarizations, *Proc. Amer. Math. Soc.* **134(1)** (2006), 177–186. DOI: 10.1090/S0002-9939-05-08325-5.
- [40] J. VAN SCHAFTINGEN AND M. WILLEM, Set transformations, symmetrizations and isoperimetric inequalities, in: *Nonlinear Analysis and Applications to Physical Sciences*, Springer-Verlag Italia, Milan, 2004, pp. 135–152.
- [41] T. WETH, Symmetry of solutions to variational problems for nonlinear elliptic equations via reflection methods, *Jahresber. Dtsch. Math.-Ver.* **112(3)** (2010), 119–158. DOI: 10.1365/s13291-010-0005-4.

Nicola Abatangelo

Dipartimento di Matematica, Alma Mater Studiorum Università di Bologna, P.zza di Porta S. Donato 5, 40126 Bologna, Italy

E-mail address: nicola.abatangelo@unibo.it

Sven Jarohs

Institut für Mathematik, Goethe-Universität Frankfurt am Main, Robert-Meyer-Str. 10, 60325 Frankfurt am Main, Germany

E-mail address: jarohs@math.uni-frankfurt.de

Received on June 17, 2022.

Accepted on February 23, 2023.